#### M1 General Physics

### Particles

#### Cross-sections

# 1 Invariant one-particle phase-space

1. Show that

$$\frac{d^3\vec{p}}{(2\pi)^3 2E(|\vec{p}|)} = \frac{d^4p}{(2\pi)^4} (2\pi)\delta(p^2 - m^2)\theta(p_0)$$
(1)

where  $E = E(|\vec{p}|) = \sqrt{\vec{p}^2 + m^2}$ , and conclude about the Lorentz invariance of this oneparticle phase-space. *Hint*: use the fact that

$$\delta(f(x)) = \sum_{i} \frac{\delta(x - x_i)}{|f'(x_i)|} \tag{2}$$

where  $x_i$  are the simple roots of f(x).

\_\_\_\_\_ Solution \_\_\_\_\_

One should just write

$$\delta(p^2 - m^2)\theta(p_0) = \delta(p_0^2 - \vec{p}^2 - m^2)\theta(p_0)$$

The equation  $p_0^2 - \vec{p}^2 - m^2 = 0$  has two roots  $p_0 = \pm E(|\vec{p}|) = \sqrt{\vec{p}^2 + m^2}$ , but the positive energy constraint  $\theta(p_0)$  selects the positive one. Since

$$\frac{d}{dp_0}(p_0^2 - \vec{p}^2 - m^2)(p_0 = E) = 2E,$$

we thus have

$$\delta(p^2 - m^2)\theta(p_0) = \frac{\delta(p_0 - E)}{2E}$$

q.e.d. The Lorentz invariance is obvious from the R.H.S of Eq. (1).

2. Write  $d^3\vec{p}$  in terms of  $p = |\vec{p}|$  (beware to this rather standard notation: p here should not be confused with the 4-momentum!) and of the elementary solid angle  $d^2\Omega$ .

\_\_\_\_\_ Solution \_\_\_\_\_

$$d^3p = p^2 \, dp \, d^2\Omega \, .$$

3. Write  $d^2\Omega$  in spherical coordinates.

$$d^2\Omega = \sin\theta d\theta d\phi$$

# 2 Phase space in the center-of-mass frame

We consider the  $2 \to 2$  process  $A(p_A) B(p_B) \to C(p_C) D(p_D)$ , where A, B, C, D are particles of mass respectively equal to  $m_A, m_B, m_C, m_D$ . Our aim is to simplify the expression of the phase space

$$d(P.S) = (2\pi)^4 \delta^{(4)} (p_A + p_B - p_C - p_D) \frac{d^3 p_C}{(2\pi)^3 2E_C} \frac{d^3 p_D}{(2\pi)^3 2E_D}$$
(3)

in the center-of-mass frame. We denote  $p_C = |\vec{p}_C|$  and  $p_D = |\vec{p}_D|$ . One may use the Mandelstam variable  $s = (p_A + p_B)^2$ . In the center-of-mass frame, we denote  $p_f^* = p_C$ .

1. Show that in the center-of-mass frame,

$$d(P.S) = \frac{1}{4\pi^2} \delta^{(3)}(\vec{p}_C + \vec{p}_D) \,\delta(E_C(p_C) + E_D(p_D) - \sqrt{s}) \,\frac{d^3 p_C}{2E_C(p_C)} \frac{d^3 p_D}{2E_D(p_D)} \,, \tag{4}$$

and give the expressions of  $E_C(p_C)$  and  $E_D(p_D)$ .

\_\_\_\_\_ Solution \_

In the center-of-mass frame, we have  $\vec{p}_A + \vec{p}_B = 0$  and thus  $s = (p_A + p_B)^2 = (E_A + E_B)^2$ . The 4-momenta conservation then reads  $\vec{p}_C + \vec{p}_D = 0$ , so that  $p_C = p_D$ , and  $E_C + E_D = E_A + E_B = \sqrt{s}$ . Besides,  $E_C = \sqrt{m_C^2 + p_C^2}$  and  $E_D = \sqrt{m_D^2 + p_D^2} = \sqrt{m_D^2 + p_C^2}$ . Thus,  $\delta^{(4)}(p_A + p_B - p_C - p_D) = \delta^{(3)}(\vec{p}_C + \vec{p}_D) \,\delta(E_C(p_C) + E_D(p_C) - \sqrt{s})$ .

2. Show finally that

$$d(P.S) = \frac{1}{4\pi^2} \frac{p_f^*}{4\sqrt{s}} d^2 \Omega \,.$$
(5)

*Hint:* one may use Eq. (2).

\_\_\_\_\_ Solution \_

We have

$$d(P.S) = \frac{1}{4\pi^2} \delta(f(p_C)) \frac{p_C^2 \, dp_C \, d^2\Omega}{2E_C(p_C) \, 2E_D(p_C)}$$

with  $f(p_C) = \sqrt{m_C^2 + p_C^2} + \sqrt{m_D^2 + p_C^2} - \sqrt{s}$ . We denote  $p_f^*$  the positive root of  $f(p_C)$ . Since

$$f'(p_C) = \frac{p_C}{\sqrt{m_C^2 + p_C^2}} + \frac{p_C}{\sqrt{m_D^2 + p_C^2}} = \frac{p_C(\sqrt{m_C^2 + p_C^2} + \sqrt{m_D^2 + p_C^2})}{\sqrt{m_C^2 + p_C^2}\sqrt{m_D^2 + p_C^2}} = \frac{p_C\sqrt{s}}{E_C(p_C)E_D(p_C)},$$

we thus get

$$\delta(f(p_C)) \ \frac{p_C^2 \, dp_C}{2E_C(p_C) \, 2E_D(p_C)} = \delta(p_C - p_f^*) \ \frac{1}{|f'(p_f^*)|} \ \frac{p_C^2 \, dp_C}{2E_C(p_C) \, 2E_D(p_C)} = \frac{p_f^*}{4\sqrt{s}}$$

and thus

$$d(P.S) = \frac{1}{4\pi^2} \frac{p_f^*}{4\sqrt{s}} d^2 \Omega \,.$$

# 3 Study of the "spinless" electron-muon scattering

Consider "spinless" electron-muon scattering. Denote  $\theta$  the scattering angle in the center-ofmass system (c.m.s), i.e. the angle between the outgoing and incoming electron (or muon) momentum. One may use the notation of section 1, with  $m_e = m_A = m_C$  and  $m_\mu = m_B = m_D$ .

1. Write the expression of the scattering amplitude  $\mathcal{M}$ .

\_\_\_\_\_ Solution \_\_\_\_\_

The scattering amplitude  $\mathcal{M}$  reads

$$i\mathcal{M} = \left[ie(p_A + p_C)^{\mu}\right] \left[-i\frac{g_{\mu\nu}}{q^2}\right] \left[ie(p_B + p_D)^{\mu}\right],$$

where  $q^2 = t = (p_A - p_C)^2$ , and thus

$$\mathcal{M} = e^2(p_A + p_C) \cdot (p_B + p_D) \frac{1}{q^2},$$

2. We denote  $s = (p_A + p_B)^2$ . In the c.m.s., write the equations satisfied by  $\vec{p}_A$ ,  $\vec{p}_B$ ,  $\vec{p}_C$ ,  $\vec{p}_D$  and  $E_A$ ,  $E_B$ ,  $E_C$ ,  $E_D$ . Deduce an equation satisfied by  $|\vec{p}_A|$  and  $|\vec{p}_C|$  and conclude about their relative magnitude.

Then, write the energy and space components of  $p_A$ ,  $p_B$ ,  $p_C$ ,  $p_D$  in terms of  $s = (p_A + p_B)^2$ and of  $E_A$ ,  $E_B$ ,  $\vec{p}_A$  and  $\vec{p}_C$ . In the c.m.s, one has  $\vec{p}_A + \vec{p}_B = 0$  and  $\vec{p}_C + \vec{p}_D = 0$ , and thus  $\sqrt{s} = E_A + E_B = E_C + E_D$ . First, this implies that

$$|\vec{p}_A| = |\vec{p}_B| = \sqrt{E_A^2 - m_A^2} = \sqrt{E_B^2 - m_B^2}$$
 and  $|\vec{p}_C| = |\vec{p}_D| = \sqrt{E_C^2 - m_A^2} = \sqrt{E_D^2 - m_B^2}$ .

Second,  $\sqrt{s} = E_A + E_B = E_C + E_D$  reads

$$\sqrt{m_A^2 + \vec{p}_A^2} + \sqrt{m_B^2 + \vec{p}_A^2} = \sqrt{m_A^2 + \vec{p}_C^2} + \sqrt{m_B^2 + \vec{p}_C^2}$$

and thus  $|\vec{p}_A| = |\vec{p}_C|$ . We can thus write

$$p_A = (E_A, \vec{p}_A)$$

$$p_B = (E_B, -\vec{p}_A)$$

$$p_C = (E_A, \vec{p}_C)$$

$$p_D = (E_B, -\vec{p}_C)$$

3. Give the expression of  $q^2$  as a function of  $\theta$  and  $|\vec{p}_A|$ . Then, write  $q^2$  in terms of  $s = (p_A + p_B)^2$ ,  $m_A$ ,  $m_B$ . One may use the obtained expression for  $|\vec{p}_A|$  in the 2020 mid-term exam, or directly solve the equation satisfied by  $|\vec{p}_A|$  in question 2.

\_ Solution \_\_

We have

$$q^{2} = (p_{A} - p_{C})^{2} = p_{A}^{2} + p_{C}^{2} - 2p_{A} \cdot p_{C} = 2m_{A}^{2} - 2E_{A}E_{C} + 2|\vec{p}_{A}||\vec{p}_{C}|\cos\theta$$
  
$$= 2m_{A}^{2} - 2E_{A}^{2} + 2|\vec{p}_{A}|^{2}\cos\theta = -2|\vec{p}_{A}|^{2}(1 - \cos\theta)$$
  
$$= -\frac{[s - (m_{A} - m_{B})^{2}][s - (m_{A} + m_{B})^{2}]}{2s}(1 - \cos\theta)$$

where we have used the fact that in the c.ms.,

$$|\vec{p}_A|^2 = \frac{[s - (m_A - m_B)^2][s - (m_A + m_B)^2]}{4s},$$

obtained either by solving

$$\sqrt{m_A^2 + \vec{p}_A^2} + \sqrt{m_B^2 + \vec{p}_A^2} = \sqrt{s}$$

or using Eq. (13) of the mid-term exam of November 2020.

4. Express the numerator of  $\mathcal{M}$  as a function of  $E_A$ ,  $E_B$ ,  $\vec{p}_A^2$  and  $\cos\theta$ . Write  $E_A$  and  $E_B$  in terms of s,  $m_A$ ,  $m_B$  (one may rely on results obtained in the 2020 mid-term exam) and show finally that

$$\mathcal{M} = e^2 \left[ \frac{3 + \cos\theta}{1 - \cos\theta} + \frac{C}{1 - \cos\theta} \right] \tag{6}$$

where C is a function of  $s, m_A, m_B$  which vanishes in the high-energy limit.

\_\_\_\_\_ Solution \_\_\_\_\_

$$(p_A + p_C) \cdot (p_B + p_D) = (2E_A, \vec{p}_A + \vec{p}_C) \cdot (2E_B, -\vec{p}_A - \vec{p}_C) = 4E_A E_B + 2\vec{p}_A^2 (1 + \cos\theta).$$

We have, either computing  $E_A = \sqrt{m_A^2 + \vec{p}_A^2}$  and  $E_B = \sqrt{m_B^2 + \vec{p}_A^2}$  from the obtained expression for  $\vec{p}_A^2$ , or using Eq. (16) of the 2020 mid-term exam,

$$E_A = \frac{s + m_A^2 - m_B^2}{2\sqrt{s}}$$
 and  $E_B = \frac{s + m_B^2 - m_A^2}{2\sqrt{s}}$ 

and thus

$$(p_A + p_C) \cdot (p_B + p_D) = \frac{2(s + m_A^2 - m_B^2)(s + m_B^2 - m_A^2) + (s - (m_A - m_B)^2)(s - (m_A + m_B)^2)(1 + \cos\theta)}{2s}$$

so that

$$\mathcal{M} = e^{2} \left[ 2 \frac{s^{2} - (m_{A} - m_{B})^{2} (m_{A} + m_{B})^{2}}{(s - (m_{A} - m_{B})^{2})(s - (m_{A} + m_{B})^{2})} \frac{1}{1 - \cos \theta} + \frac{1 + \cos \theta}{1 - \cos \theta} \right]$$
$$= e^{2} \left[ \frac{3 + \cos \theta}{1 - \cos \theta} + 4 \frac{s(m_{A}^{2} + m_{B}^{2}) - (m_{A}^{2} - m_{B}^{2})}{(s - (m_{A} - m_{B})^{2})(s - (m_{A} + m_{B})^{2})} \frac{1}{1 - \cos \theta} \right]$$

#### 5. Prove finally that the differential cross-section reads

$$\left. \frac{d\sigma}{d\Omega} \right|_{cms} = \frac{\alpha^2}{4s} \left( \frac{3 + C + \cos\theta}{1 - \cos\theta} \right)^2,\tag{7}$$

where  $\alpha = e^2/(4\pi)$  is the fine-structure constant.

\_\_\_\_\_ Solution \_\_\_

We know that in the c.ms.,  $p_i^* = |\vec{p}_A| = |\vec{p}_C| = p_f^*$  so that the differential cross-section reads

$$\frac{d\sigma}{d\Omega}\Big|_{cms} = \frac{1}{64\pi^2 s} \frac{p_f^*}{p_i^*} |\mathcal{M}|^2 = \frac{e^4}{64\pi^2 s} \left(\frac{3+C+\cos\theta}{1-\cos\theta}\right)^2 = \frac{\alpha^2}{4s} \left(\frac{3+C+\cos\theta}{1-\cos\theta}\right)^2.$$