

Electromagnetic mass splittings in a SU(3) octet

Note : in Mathematica, greek letters are obtained by using the key "Esc" before and after typing the abbreviation of the greek letter with a latine spelling

Theory

■ octet content and properties

■ octet of mesons $J^P = 0^-$

$$\Phi = \{\{1 / \text{Sqrt}[2] \pi^0 - 1 / \text{Sqrt}[6] \eta, \pi^p, K^p\}, \\ \{\pi^m, -1 / \text{Sqrt}[2] \pi^0 - 1 / \text{Sqrt}[6] \eta, K^0\}, \{K^m, K^{0b}, \text{Sqrt}[2/3] \eta\}\}$$

$$\begin{pmatrix} \frac{\pi^0}{\sqrt{2}} - \frac{\eta}{\sqrt{6}} & \pi^p & K^p \\ \pi^m & -\frac{\eta}{\sqrt{6}} - \frac{\pi^0}{\sqrt{2}} & K^0 \\ K^m & K^{0b} & \sqrt{\frac{2}{3}} \eta \end{pmatrix}$$

■ octet of baryons $J^P = \frac{1}{2}^+$

$$\Psi = \{\{1 / \text{Sqrt}[2] \Sigma^0 - 1 / \text{Sqrt}[6] \Lambda, \Sigma^p, p\}, \\ \{\Sigma^m, -1 / \text{Sqrt}[2] \Sigma^0 - 1 / \text{Sqrt}[6] \Lambda, n\}, \{\Xi^m, \Xi^0, \text{Sqrt}[2/3] \Lambda\}\}$$

$$\begin{pmatrix} \frac{\Sigma^0}{\sqrt{2}} - \frac{\Lambda}{\sqrt{6}} & \Sigma^p & p \\ \Sigma^m & -\frac{\Lambda}{\sqrt{6}} - \frac{\Sigma^0}{\sqrt{2}} & n \\ \Xi^m & \Xi^0 & \sqrt{\frac{2}{3}} \Lambda \end{pmatrix}$$

One should be careful: Σ and Σ' are not charge conjugated

$$\Psi_b = \{\{-(\Lambda b / \text{Sqrt}[6]) + \Sigma^{0b} / \text{Sqrt}[2], \Sigma^{mb}, \Sigma^p\}, \\ \{\Sigma^{pb}, -(\Lambda b / \text{Sqrt}[6]) - \Sigma^{0b} / \text{Sqrt}[2], \Sigma^{0b}\}, \{pb, nb, \text{Sqrt}[2/3] * \Lambda b\}\}$$

$$\begin{pmatrix} \frac{\Sigma^{0b}}{\sqrt{2}} - \frac{\Lambda b}{\sqrt{6}} & \Sigma^{mb} & \Sigma^p \\ \Sigma^{pb} & -\frac{\Lambda b}{\sqrt{6}} - \frac{\Sigma^{0b}}{\sqrt{2}} & \Sigma^{0b} \\ pb & nb & \sqrt{\frac{2}{3}} \Lambda b \end{pmatrix}$$

the isospin, hypercharge and charge operators

$$\mathbf{Iz} = 1/2 \{\{1, 0, 0\}, \{0, -1, 0\}, \{0, 0, 0\}\}$$

$$\begin{pmatrix} \frac{1}{2} & 0 & 0 \\ 0 & -\frac{1}{2} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\mathbf{Y} = 1/3 \{\{1, 0, 0\}, \{0, 1, 0\}, \{0, 0, -2\}\}$$

$$\begin{pmatrix} \frac{1}{3} & 0 & 0 \\ 0 & \frac{1}{3} & 0 \\ 0 & 0 & -\frac{2}{3} \end{pmatrix}$$

$$\mathbf{Q} = 1/2 \mathbf{Y} + \mathbf{Iz}$$

$$\begin{pmatrix} \frac{2}{3} & 0 & 0 \\ 0 & -\frac{1}{3} & 0 \\ 0 & 0 & -\frac{1}{3} \end{pmatrix}$$

■ the isospin, hypercharge and charge of the octet of mesons $J^P = 0^-$

η is isoscalar, (K^0, K^+) and (K^-, K^0_b) are two isodoublet, (π^-, π^0, π^+) is an isovector:

$$\text{Simplify}[\mathbf{Iz}.\Phi - \Phi.\mathbf{Iz}]$$

$$\begin{pmatrix} 0 & \pi p & \frac{K_p}{2} \\ -\pi m & 0 & -\frac{K_0}{2} \\ -\frac{K_m}{2} & \frac{K_0 b}{2} & 0 \end{pmatrix}$$

$$\text{Simplify}[\mathbf{Y}.\Phi - \Phi.\mathbf{Y}]$$

$$\begin{pmatrix} 0 & 0 & \frac{K_p}{2} \\ 0 & 0 & K_0 \\ -K_m & -K_0 b & 0 \end{pmatrix}$$

$$\text{Simplify}[\mathbf{Q}.\Phi - \Phi.\mathbf{Q}]$$

$$\begin{pmatrix} 0 & \pi p & K_p \\ -\pi m & 0 & 0 \\ -K_m & 0 & 0 \end{pmatrix}$$

■ the isospin, hypercharge and charge of octet of baryons $J^P = \frac{1}{2}^+$

Λ^0 is an isoscalar, (n, p) and (Ξ^-, Ξ^0) are two isodoublet, $(\Sigma^-, \Sigma^0, \Sigma^+)$ is an isovector:

```
Simplify[Iz. $\Psi$  -  $\Psi$ .Iz]
```

$$\begin{pmatrix} 0 & \Sigma p & \frac{p}{2} \\ -\Sigma m & 0 & -\frac{n}{2} \\ -\frac{\Xi m}{2} & \frac{\Xi 0}{2} & 0 \end{pmatrix}$$

```
Simplify[Y. $\Psi$  -  $\Psi$ .Y]
```

$$\begin{pmatrix} 0 & 0 & p \\ 0 & 0 & n \\ -\Xi m & -\Xi 0 & 0 \end{pmatrix}$$

```
Simplify[Q. $\Psi$  -  $\Psi$ .Q]
```

$$\begin{pmatrix} 0 & \Sigma p & p \\ -\Sigma m & 0 & 0 \\ -\Xi m & 0 & 0 \end{pmatrix}$$

■ Mass splittings

■ the four invariant operators

```
O1 =  $\Psi$ b.Q.Q. $\Psi$ ;
```

```
O2 =  $\Psi$ b. $\Psi$ .Q.Q;
```

```
O3 =  $\Psi$ b.Q. $\Psi$ .Q;
```

```
O4 =  $\Psi$ b. $\Psi$ ;
```

```
O $\alpha$  = Expand[Tr[O1]]
```

$$\frac{n nb}{9} + \frac{4 p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} + \frac{\Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} + \frac{\Sigma m \Sigma m b}{9} + \frac{4 \Sigma p \Sigma p b}{9}$$

```
O $\beta$  = Expand[Tr[O2]]
```

$$\frac{n nb}{9} + \frac{p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} + \frac{4 \Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} + \frac{4 \Sigma m \Sigma m b}{9} + \frac{\Sigma p \Sigma p b}{9}$$

```
O $\gamma$  = Expand[Tr[O3]]
```

$$\frac{n nb}{9} - \frac{2 p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} - \frac{2 \Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} - \frac{2 \Sigma m \Sigma m b}{9} - \frac{2 \Sigma p \Sigma p b}{9}$$

Tr [Ψ Ψ b] is the identity operator:

```
Oid = Expand[Tr[O4]]
```

$$n nb + p pb + \Lambda \Lambda b + \Xi 0 \Xi 0 b + \Xi m \Xi p + \Sigma 0 \Sigma 0 b + \Sigma m \Sigma m b + \Sigma p \Sigma p b$$

■ the relevant mass operator

Note: since Oid is proportional to Id, it does not contribute to mass differences ; we anyway expand the mass operator as an expansion on the 4 invariants

$$M = \alpha O\alpha + \beta O\beta + \gamma O\gamma + \delta Oid$$

$$\begin{aligned} & \gamma \left(\frac{n nb}{9} - \frac{2 p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} - \frac{2 \Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} - \frac{2 \Sigma m \Sigma m b}{9} - \frac{2 \Sigma p \Sigma p b}{9} \right) + \\ & \beta \left(\frac{n nb}{9} + \frac{p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} + \frac{4 \Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} + \frac{4 \Sigma m \Sigma m b}{9} + \frac{\Sigma p \Sigma p b}{9} \right) + \\ & \alpha \left(\frac{n nb}{9} + \frac{4 p pb}{9} + \frac{\Lambda \Lambda b}{6} + \frac{\Xi 0 \Xi 0 b}{9} + \frac{\Xi m \Xi p}{9} - \frac{\Lambda b \Sigma 0}{6 \sqrt{3}} - \frac{\Lambda \Sigma 0 b}{6 \sqrt{3}} + \frac{5 \Sigma 0 \Sigma 0 b}{18} + \frac{\Sigma m \Sigma m b}{9} + \frac{4 \Sigma p \Sigma p b}{9} \right) + \\ & \delta (n nb + p pb + \Lambda \Lambda b + \Xi 0 \Xi 0 b + \Xi m \Xi p + \Sigma 0 \Sigma 0 b + \Sigma m \Sigma m b + \Sigma p \Sigma p b) \end{aligned}$$

Mp = Coefficient[M, p pb]

$$\frac{4 \alpha}{9} + \frac{\beta}{9} - \frac{2 \gamma}{9} + \delta$$

Mn = Coefficient[M, n nb]

$$\frac{\alpha}{9} + \frac{\beta}{9} + \frac{\gamma}{9} + \delta$$

MΛ = Coefficient[M, Λ Λb]

$$\frac{\alpha}{6} + \frac{\beta}{6} + \frac{\gamma}{6} + \delta$$

MΣp = Coefficient[M, Σp Σpb]

$$\frac{4 \alpha}{9} + \frac{\beta}{9} - \frac{2 \gamma}{9} + \delta$$

MΣ0 = Coefficient[M, Σ0 Σ0b]

$$\frac{5 \alpha}{18} + \frac{5 \beta}{18} + \frac{5 \gamma}{18} + \delta$$

MΣm = Coefficient[M, Σm Σmb]

$$\frac{\alpha}{9} + \frac{4 \beta}{9} - \frac{2 \gamma}{9} + \delta$$

MΞ0 = Coefficient[M, Ξ0 Ξ0b]

$$\frac{\alpha}{9} + \frac{\beta}{9} + \frac{\gamma}{9} + \delta$$

MΞm = Coefficient[M, Ξm Ξp]

$$\frac{\alpha}{9} + \frac{4 \beta}{9} - \frac{2 \gamma}{9} + \delta$$

{MΞm - MΞ0, MΣm - MΣp + Mp - Mn}

$$\left\{ \frac{\beta}{3} - \frac{\gamma}{3}, \frac{\beta}{3} - \frac{\gamma}{3} \right\}$$

Conclusion : MΞ m- MΞ 0= MΣ m- MΣ p+ Mp - Mn

Phenomenology

- Experimental data for masses (in MeV/c²) of the octet of baryons $J^P = \frac{1}{2}^+$

$M_{nE} = 939.56$;
 $M_{pE} = 938.27$;
 $M_{\Xi mE} = 1321.71$;
 $M_{\Xi 0E} = 1314.86$;
 $M_{\Sigma mE} = 1197.45$;
 $M_{\Sigma 0E} = 1192.64$; $M_{\Sigma pE} = 1189.37$;

- Comparison with theory

$$\{M_{\Xi mE} - M_{\Xi 0E}, M_{\Sigma mE} - M_{\Sigma pE} + M_{pE} - M_{nE}\}$$

$$\{6.85, 6.79\}$$

$$(M_{\Xi mE} - M_{\Xi 0E}) / (M_{\Sigma mE} - M_{\Sigma pE} + M_{pE} - M_{nE})$$

$$1.00884$$