

q -deformed Painlevé τ function and q -deformed conformal blocks

M. Bershtein

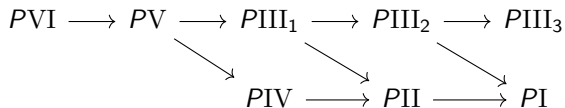
based on joint papers with A. Shchekkin
arXiv:1608.02566 and 1608.02568

20 October 2016

Painlevé equations

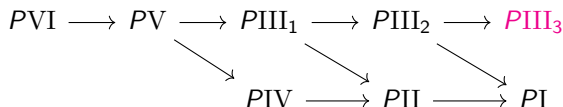
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Each arrow stands for degeneration of the equation. Painlevé VI depend on 4 parameters, Painlevé V depend on 3 parameters, ..., Painlevé I and Painlevé III₃ do not have parameters.

- Today we concentrate on Painlevé III₃

$$w'' = \frac{w'^2}{w} - \frac{w'}{z} + \frac{2w^2}{z^2} - \frac{2}{z}$$

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- Painleve III₃ corresponds to $D_8^{(1)}$ surface and has $D_0^{(1)}$ symmetry.

Painlevé III $D_8^{(1)}$ equation

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- Second order nonlinear differential equation

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where $D_{[\log z]}^2$ denotes second Hirota operator with respect to $\log z$.

The function $w(z)$ is equal to $-z^{1/2} \tau(z)^2 / \tau_1(z)^2$.

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- Algebraic solution $w = \pm\sqrt{z}, \tau(z) = z^{1/16} e^{\pm 4z^{1/2}}$.

This solution is invariant under Bäcklund transformation π .

Gamayun-Iorgov-Lisovyy formula

- [Gamayun-Iorgov-Lisovyy 11]

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- Algebraic solution corresponds to $\sigma = 1/4, s = \pm 1$. We have an identity

$$z^{1/16} e^{\mp 4\sqrt{z}} = \sum_{n \in \mathbb{Z}} (\mp 1)^n B_n z^{n^2+n/2} \mathcal{F}((1/4+n)^2|z),$$

where B_n are certain rational numbers. Values $\Delta = (n+1/4)^2$ are conformal dimensions of twist fields [Al. Zamolodchikov 87], [Gavrylenko, Marshakov 15].

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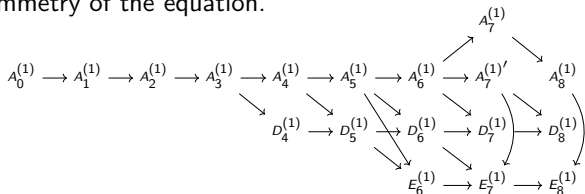
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Geometric approach to difference Painlevé equations [Sakai 01]: for each equation one can assign two lattices, R which encodes phase space of the equations and R^\perp which encodes symmetry of the equation.

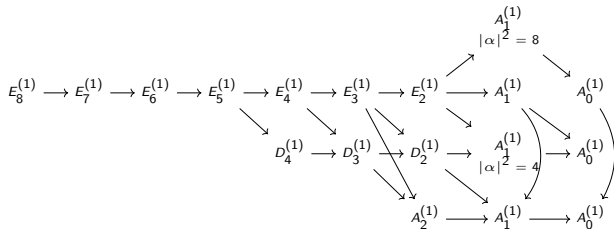
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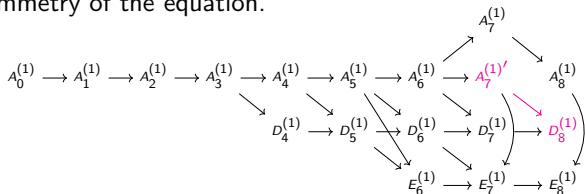
Lattice R^\perp



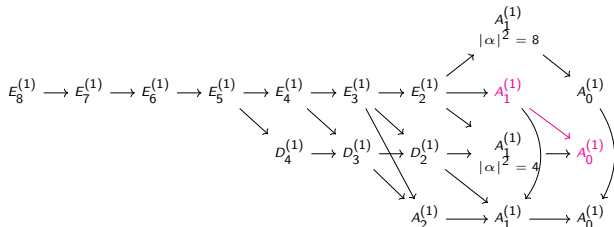
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Lattice R



Lattice R^\perp



We consider q -difference equation with $R = A_7^{(1)'}$ and $R^\perp = A_1^{(1)}$.

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- Moduli space of such surfaces X is parametrized by coordinates Z, q, F, G denotes coordinates on the surface X .
- W is equal to $Dih_4 \times W(A_1^{(1)})$ where Dih_4 is dihedral group of square Group W can be presented by generators s_0, s_1, π_1, π_2 and relations as

$$s_0^2 = s_1^2 = \pi_1^2 = \pi_2^2 = (\pi_1\pi_2)^2 = 1, \quad s_1 = \pi_2 s_0 \pi_2^{-1}, \quad s_0 = \pi_2^2 s_0 \pi_2^{-2} = \pi_1 s_0 \pi_1^{-1},$$

where s_0, s_1 are Weyl group simple reflections, π_1, π_2 are generators of Dih_4 (symmetry and rotation on 90°).

- The action of W is given in the table

	Z	q	F	G
π_1	Z^{-1}	q^{-1}	$q^{-1}Z^{-1}F$	G^{-1}
π_2	$Z^{-1}q^{-1}$	q	$Z^{-1}G$	F^{-1}
s_1	Z^{-1}	q	$F \frac{(G-1)^2}{(G-Z)^2}$	$Z^{-1}G$
s_0	$Z^{-1}q^{-2}$	q	$q^{-1}Z^{-1}F$	$G \frac{(1-F)^2}{(Zq-F)^2}$

Discrete dynamics

- $T = \pi_2^{-1} \circ s_0$, element of infinite order (translation) in W .

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$$(\bar{Z}, \bar{q}, \bar{F}, \bar{G}) = \left(qZ, q, \frac{(F - qZ)^2}{(F - 1)^2 G}, F \right),$$
$$(\underline{Z}, \underline{q}, \underline{F}, \underline{G}) = \left(q^{-1}Z, q, G, \frac{(G - Z)^2}{F(G - 1)^2} \right).$$

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- Therefore

$$\bar{G}\underline{G} = \left(\frac{G - Z}{G - 1} \right)^2.$$

In the limit $q \rightarrow 1$ this equation goes to Painlevé III₃($D_8^{(1)}$) equation.

τ functions

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Theorem

Action of generators s_1, π_1, π_2 of group W on $\mathcal{T}_i, i = \overline{1, 4}$ given below provides a representation of W in the field $\mathbb{C}(\mathcal{T}_1, \mathcal{T}_2, \mathcal{T}_3, \mathcal{T}_4, q^{1/4}, Z^{1/4})$.

	Z	q	F	G	\mathcal{T}_1	\mathcal{T}_2	\mathcal{T}_3	\mathcal{T}_4
π_1	$1/Z$	$1/q$	$F/(qZ)$	$1/G$	\mathcal{T}_3	\mathcal{T}_2	\mathcal{T}_1	\mathcal{T}_4
π_2	$1/(qZ)$	q	G/Z	$1/F$	\mathcal{T}_4	\mathcal{T}_1	\mathcal{T}_2	\mathcal{T}_3
s_1	$1/Z$	q	$F \frac{(G-1)^2}{(G-Z)^2}$	G/Z	\mathcal{T}_1	$\frac{\mathcal{T}_3^2 + Z^{1/2} \mathcal{T}_1^2}{Z^{1/4} \mathcal{T}_4}$	\mathcal{T}_3	$\frac{\mathcal{T}_1^2 + Z^{1/2} \mathcal{T}_3^2}{Z^{1/4} \mathcal{T}_2}$
s_0	$1/(q^2 Z)$	q	$F/(qZ)$	$G \frac{(1-F)^2}{(Zq-F)^2}$	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4
T	qZ	q	$\frac{(F-qZ)^2}{(-1+F)^2 G}$	F	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$

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s_1	$1/Z$	q	$F \frac{(G-1)^2}{(G-Z)^2}$	G/Z	\mathcal{T}_1	$\frac{\mathcal{T}_3^2 + Z^{1/2} \mathcal{T}_1^2}{Z^{1/4} \mathcal{T}_4}$	\mathcal{T}_3	$\frac{\mathcal{T}_1^2 + Z^{1/2} \mathcal{T}_3^2}{Z^{1/4} \mathcal{T}_2}$
s_0	$1/(q^2 Z)$	q	$F/(qZ)$	$G \frac{(1-F)^2}{(Zq-F)^2}$	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4
T	qZ	q	$\frac{(F-qZ)^2}{(-1+F)^2 G}$	F	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$

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τ functions

- We follow approach [Tsuda 06] (cf. [Noumi 16])

Theorem

Action of generators s_1, π_1, π_2 of group W on $\mathcal{T}_i, i = \overline{1, 4}$ given below provides a representation of W in the field $\mathbb{C}(\mathcal{T}_1, \mathcal{T}_2, \mathcal{T}_3, \mathcal{T}_4, q^{1/4}, Z^{1/4})$.

	Z	q	F	G	\mathcal{T}_1	\mathcal{T}_2	\mathcal{T}_3	\mathcal{T}_4
π_1	$1/Z$	$1/q$	$F/(qZ)$	$1/G$	\mathcal{T}_3	\mathcal{T}_2	\mathcal{T}_1	\mathcal{T}_4
π_2	$1/(qZ)$	q	G/Z	$1/F$	\mathcal{T}_4	\mathcal{T}_1	\mathcal{T}_2	\mathcal{T}_3
s_1	$1/Z$	q	$F \frac{(G-1)^2}{(G-Z)^2}$	G/Z	\mathcal{T}_1	$\frac{\mathcal{T}_3^2 + Z^{1/2} \mathcal{T}_1^2}{Z^{1/4} \mathcal{T}_4}$	\mathcal{T}_3	$\frac{\mathcal{T}_1^2 + Z^{1/2} \mathcal{T}_3^2}{Z^{1/4} \mathcal{T}_2}$
s_0	$1/(q^2 Z)$	q	$F/(qZ)$	$G \frac{(1-F)^2}{(Zq-F)^2}$	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4
T	qZ	q	$\frac{(F-qZ)^2}{(-1+F)^2 G}$	F	\mathcal{T}_2	$\frac{\mathcal{T}_2^2 + (qZ)^{1/2} \mathcal{T}_4^2}{(qZ)^{1/4} \mathcal{T}_1}$	\mathcal{T}_4	$\frac{\mathcal{T}_4^2 + (qZ)^{1/2} \mathcal{T}_2^2}{(qZ)^{1/4} \mathcal{T}_3}$

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s_1	$1/Z$	q	$F \frac{(G-1)^2}{(G-Z)^2}$	G/Z	\mathcal{T}_1	$\frac{\mathcal{T}_3^2 + Z^{1/2} \mathcal{T}_1^2}{Z^{1/4} \mathcal{T}_4}$	\mathcal{T}_3	$\frac{\mathcal{T}_1^2 + Z^{1/2} \mathcal{T}_3^2}{Z^{1/4} \mathcal{T}_2}$
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- Action W can be expressed as composition of mutations in cluster algebra. Therefore it maps \mathcal{T}_i to *Laurent polynomials*.

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Let $|q| \neq 1$ and $u \neq q^n$, $n \in \mathbb{Z}$. Then series (2) converges uniformly and absolutely on every bounded subset of \mathbb{C} .

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- $\mathcal{F}(u; q^{-1}, q|Z)/(uq; q, q)_\infty^2$ is topological string partition function for local $\mathbb{P}^1 \times \mathbb{P}^1$ geometry.

Definition

Function $\mathcal{T}(u, s; q|Z)$ given by the formula

$$\mathcal{T}(u, s; q|Z) = R(u; q|Z) \sum_{n \in \mathbb{Z}} s^n C(uq^{2n}; q|Z) \frac{\mathcal{F}(uq^{2n}; q^{-1}, q|Z)}{(uq^{2n+1}; q, q)_{\infty} (u^{-1}q^{-2n-1}; q, q)_{\infty}}$$

is called q -deformed τ function of Painlevé III(D_8) equation if function $C(u; q|Z)$ satisfy equations

$$C(uq; q|Z)C(uq^{-1}; q|Z)/C(u; q|Z)^2 = -Z^{1/2}$$

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and function $R(u; q|Z)$ satisfy homogeneous version of these equations (with r.h.s equals 1) and is q^2 -periodic in u , i.e. $R(uq^2; q|Z)/R(u; q|Z) = 1$.

If $C(u; q|Z) = C(u^{-1}; q|Z)$ and $R(u; q|Z) = R(u^{-1}; q|Z)$ then functions C, R, \mathcal{T} are called u -inverse invariant.

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Cf. [Gamayun-Iorgov-Lisovyy 11]

$$\tau(\sigma, s|z) = \sum_{n \in \mathbb{Z}} s^n z^{(\sigma+n)^2} C(\sigma+n) \mathcal{F}((\sigma+n)^2|z)$$

Properties

- The following examples of $C(u; q|Z)$ satisfy definition

$$C_1(u; q|Z) = \frac{\Gamma((qZ)^{1/4}; q^{1/4}, q^{1/4})^3}{(\Gamma(i(qZu)^{1/4}; q^{1/4}, q^{1/4}) \Gamma(i(qZ)^{1/4}u^{-1/4}; q^{1/4}, q^{1/4}))}$$

$$C_c(u; q|Z) = (-1)^{2\left(\frac{\log u}{2 \log q}\right)^2} \Gamma(-(qZ)^{1/4}; q^{1/4}, q^{1/4}) \exp\left(\frac{\log^2 u \log Z}{4 \log^2 q}\right).$$

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Conjecture

Function $\mathcal{T}(u, s; q|Z)$ satisfies bilinear equations

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Theorem

This conjecture is equivalent to

$$\begin{aligned} & \sum_{2n \in \mathbb{Z}} \frac{u^{2n} Z^{2n^2}}{\prod_{\epsilon, \epsilon' = \pm 1} (u^\epsilon q^{1+2\epsilon'n}; q, q)_\infty} \mathcal{F}(uq^{-2n}; q^{-1}, q|q^{-1}Z) \mathcal{F}(uq^{2n}; q^{-1}, q|qZ) = \\ & = (1 - Z^{1/2}) \sum_{2n \in \mathbb{Z}} \frac{Z^{2n^2}}{\prod_{\epsilon, \epsilon' = \pm 1} (u^\epsilon q^{1+2\epsilon'n}; q, q)_\infty} \mathcal{F}(uq^{-2n}; q^{-1}, q|Z) \mathcal{F}(uq^{2n}; q^{-1}, q|Z) \end{aligned}$$

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- This conjecture was checked by computer calculation up to Z^3 .
- Continuous limit of this relation gives known relation for Virasoro conformal blocks

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It follows from Conjecture that for any u -inverse invariant τ function $T(u, s; q|Z)$

$$\mathcal{T}(q^{1/2}, \pm 1; q|Z) = \frac{R(q^{1/2}|Z)C(q^{1/2}|Z)}{(q^{3/2}; q, q)_{\infty}(q^{1/2}; q, q)_{\infty}} (\mp Z^{1/2} q^{1/2}; q^{1/2}, q^{1/2})_{\infty}.$$

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$$\mathcal{T}(q^{1/2}, \pm 1; q|Z) = \frac{R(q^{1/2}|Z)C(q^{1/2}|Z)}{(q^{3/2}; q, q)_{\infty}(q^{1/2}; q, q)_{\infty}} (\mp Z^{1/2} q^{1/2}; q^{1/2}, q^{1/2})_{\infty}.$$

- This formula is equivalent to relation

$$(\mp Z^{1/2} q^{1/2}; q^{1/2}, q^{1/2})_{\infty} = \sum_{n \in \mathbb{Z}} (\mp 1)^n Z^{n^2 + n/2} P_n(q) \mathcal{F}(q^{2n+1/2}, q, q|Z),$$

Algebraic solution

- q -difference equation $G(qZ)G(q^{-1}Z) = (G(Z) - Z)^2 / (G(Z) - 1)^2$ has algebraic solution $G(Z) = \pm\sqrt{Z}$.

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- Nekrasov partition function for pure $U(1)$ 4d theory $Z(z) = \exp(z)$.
Algebraic solution for Painlevé $D_8^{(1)}$ $\tau(z) \sim z^{1/16} \exp(4\sqrt{z})$,
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Algebraic solution for q -Painlevé $A_7^{(1')}$ $\mathcal{T}(z) \sim (\mp Z^{1/2} q^{1/2}; q^{1/2}, q^{1/2})_{\infty}$,

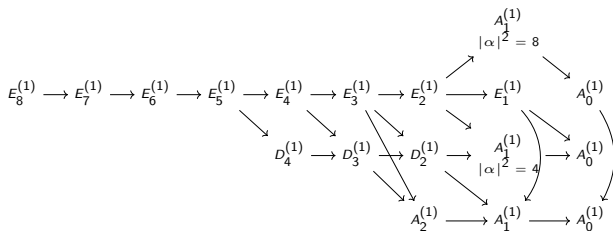
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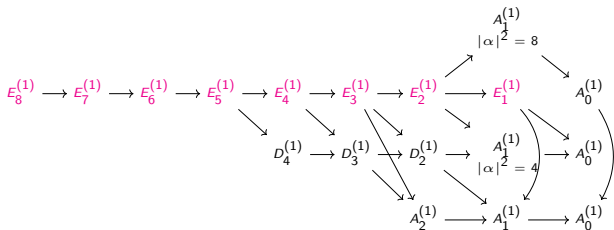
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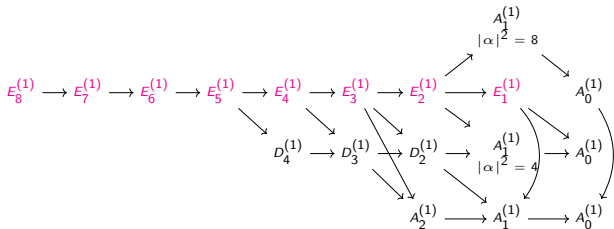
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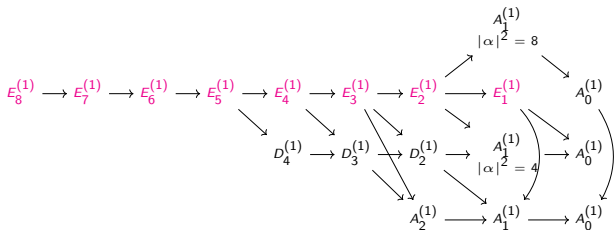
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- It is natural to conjecture that q -Painlevé equations with the symmetry type $E_{N+1}^{(1)}$ correspond to 5d Nekrasov partition functions with N fundamental multiplets.
- for $N \leq 4$ $q \rightarrow 1$ limit this relation is known [Gamayun Iorgov Lisovyy 13].
- [Seiberg 96] argued that these gauge theories has E_{N+1} global symmetry. It would be interesting to find a physical interpretation of the affine Weyl groups $E_{N+1}^{(1)}$, the symmetry group of the corresponding discrete Painlevé equations.

Thank you for your attention!

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- Elements $T, \pi_2^2 \in W$ acts on the τ -function $\mathcal{T}(u, s; q|Z)$ in a clear manner

$$T: \mathcal{T}(u, s; q|Z) \mapsto \mathcal{T}(u, s; q|qZ), \quad \text{and} \quad \pi_2^2: \mathcal{T}(u, s; q|Z) \mapsto \mathcal{T}(uq, s; q|Z)$$

It is natural to ask for the action of whole group W . The remaining transformations are $Z \mapsto Z^{-1}$ and $q \mapsto q^{-1}$. The second transformation is transparent due to $\mathcal{F}(u; q^{-1}, q|Z) = \mathcal{F}(u; q, q^{-1}|Z)$.

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- For topological string partition function for local $\mathbb{P}^1 \times \mathbb{P}^1$ geometry there exists fiber-base duality which interchange two factors \mathbb{P}^1 . In terms of the functions $\mathcal{F}(u; q^{-1}, q|Z^{-1})$ this duality has the form

$$\frac{\mathcal{F}(u; q^{-1}, q|Z)}{(uq; q, q)_{\infty}^2} = \frac{\mathcal{F}(uZ; q^{-1}, q|Z^{-1})}{(uZq; q, q)_{\infty}^2}.$$

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- [Bonelli Grassi Tanzini 16] (using [Grassi Hatsuda Marino 14]) also constructed function, which in the limit $q \rightarrow 1$ goes to the $\tau(u, s|z)$. It is interesting to note that they work in different region of q , namely $|q| = 1$ in their paper. There exists a relation between our papers.