

Reducing the combinatorial uncertainties using kinematic variables

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27 October 2011



In collaboration with K. Choi and D. Guadagnoli, arXiv:1109.2201, submitted to JHEP.

Combinatorics?

- * ... is a branch of mathematics concerning the study of finite or countable discrete structures. Aspects of combinatorics include **counting the structures of a given kind and size**, deciding when certain criteria can be met, and constructing and analyzing objects **meeting the criteria**, finding "largest", "smallest", or "optimal" objects, ...
- * A mathematician who studies combinatorics is called a *combinatorialist*.

...Wikipedia, rev. 12 Oct. 2011

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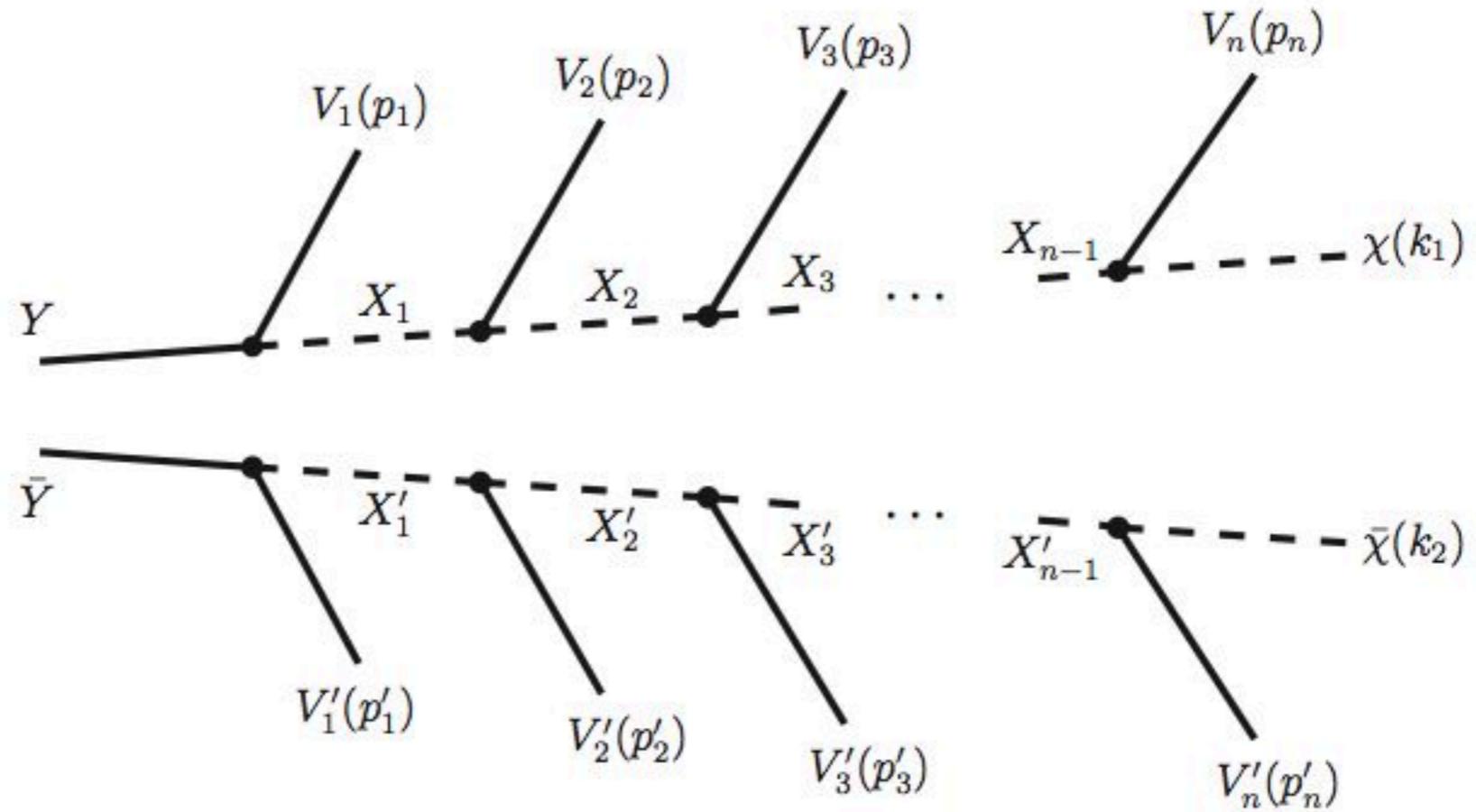
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particle physicist studying LHC

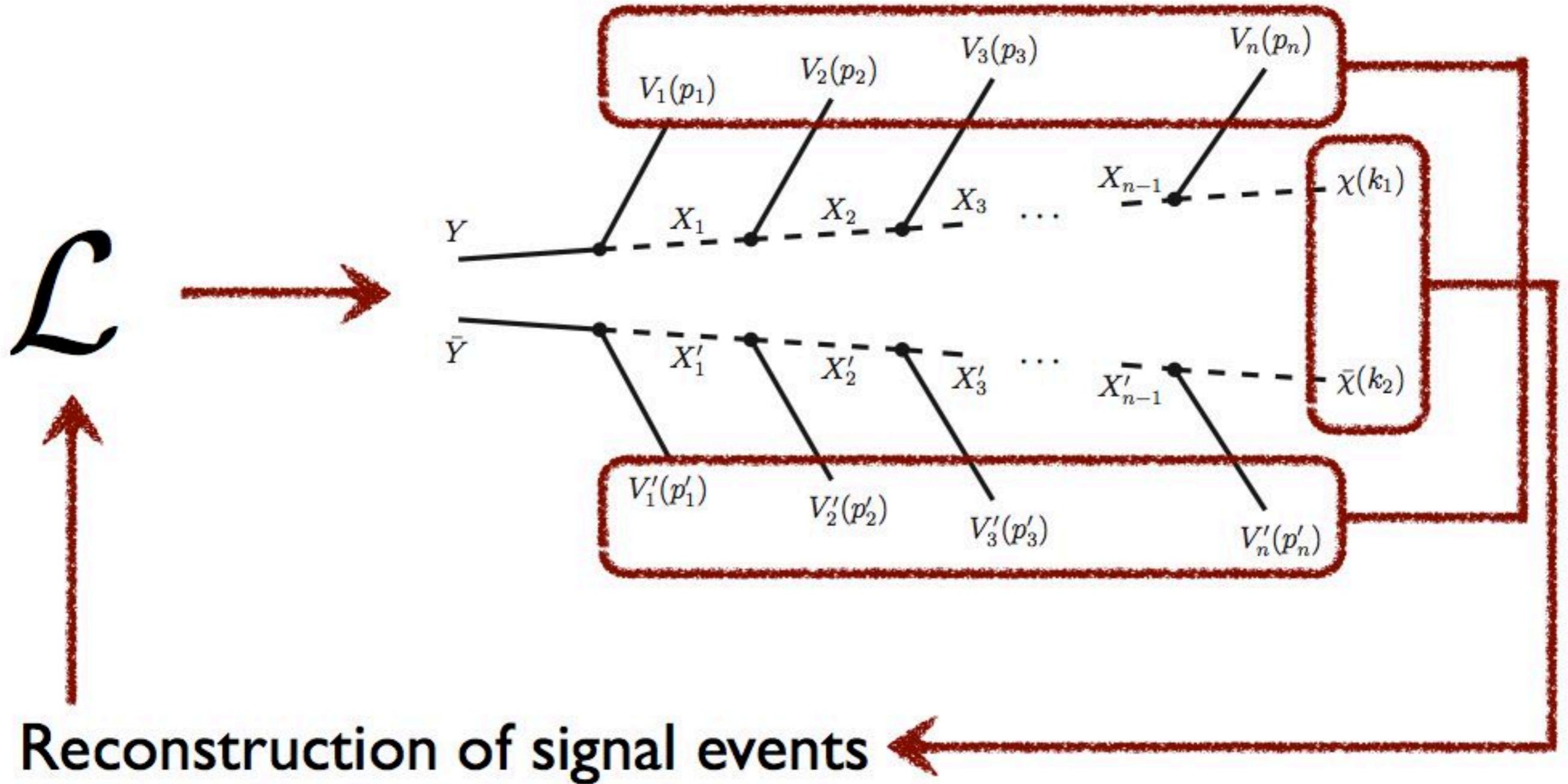


Combinatorics in collider signature

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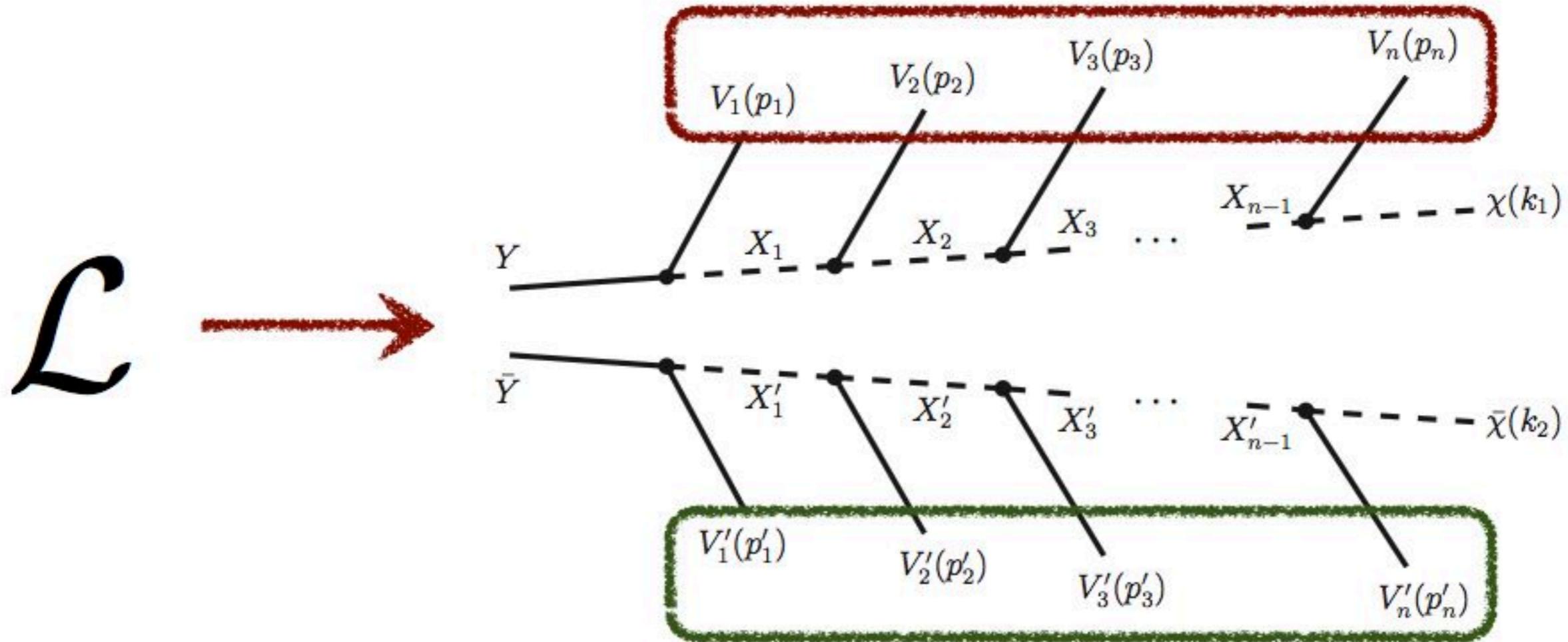


Combinatorics in collider signature



Reconstruction of signal events

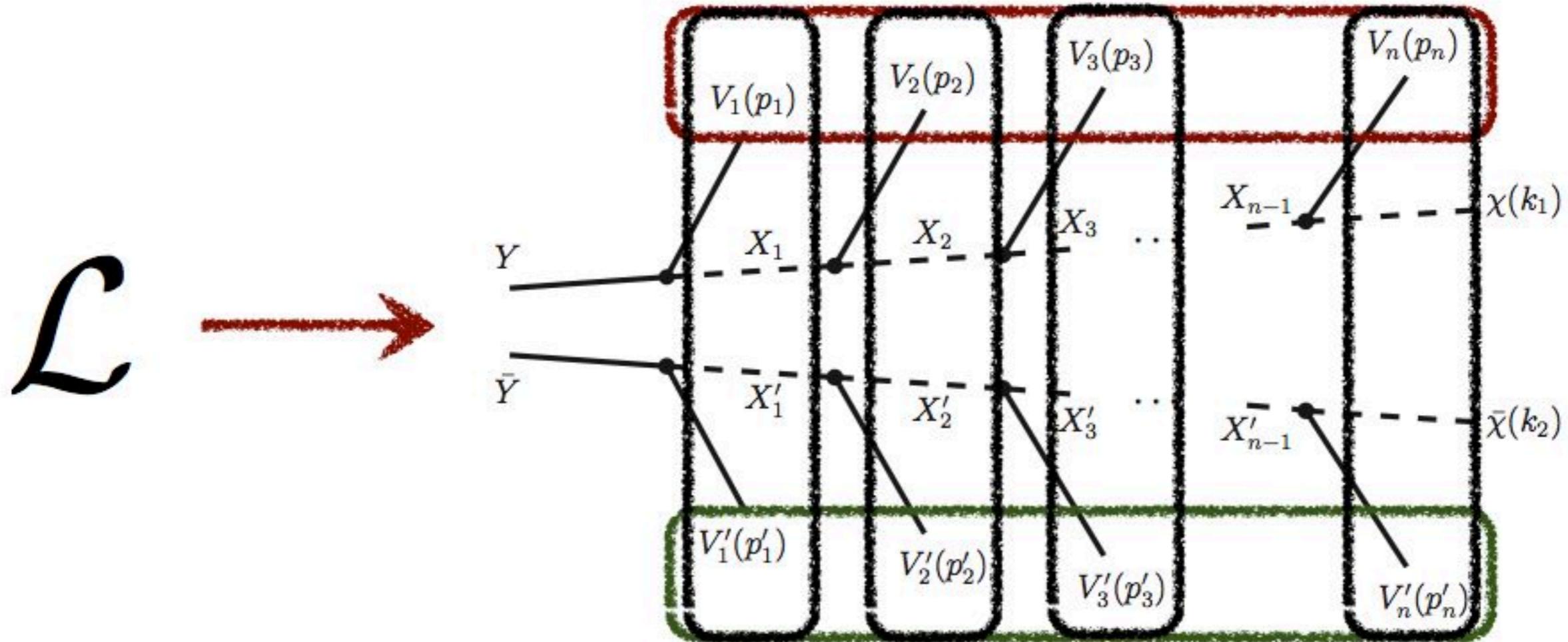
Combinatorics in collider signature



Reconstruction of signal events

* Combinatorial ambiguities of pairing

Combinatorics in collider signature



Reconstruction of signal events

* Combinatorial ambiguities of pairing and ordering

Combinatorics in collider signature: example (I)

$$Z \rightarrow e^+ e^-$$

$$Z \rightarrow \mu^+ \mu^-$$

No combinatorial ambiguity

Combinatorics in collider signature: example (I)

$$Z \rightarrow e^+ e^-$$

$$Z \rightarrow e^+ e^-$$

of possible pairings = 2

can be resolved by constructing invariant masses:

$$|(p^{e^+} + p^{e^-})^2 - m_Z^2| = 0 \quad \text{for right pairing,}$$

$$|(p^{e^+} + p^{e^-})^2 - m_Z^2| \neq 0 \quad \text{for wrong pairing,}$$

on an event-by-event basis.

Combinatorics in collider signature: example (II)

$$\tilde{q} \rightarrow \tilde{\chi}_2^0 q \rightarrow l^\pm \tilde{l}^\mp q \rightarrow \tilde{\chi}_1^0 l^\pm l^\mp q$$

$$\tilde{q} \rightarrow \tilde{\chi}_1^0 q$$

$$(m_{l^\pm l^\mp}^{\max})^2 = \frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}}^2)(m_{\tilde{l}}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}}^2}$$

Combinatorics in collider signature: example (II)

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$$\tilde{q} \rightarrow \tilde{\chi}_1^0 q$$

main background: 2 jets + 2 leptons + missing ET

$$\tilde{q} \rightarrow \tilde{\chi}_1^\pm q' \rightarrow \tilde{\chi}_1^0 l^\pm \nu q'$$

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↓ same flavor leptons

main background: 2 jets + **2 leptons** + missing ET

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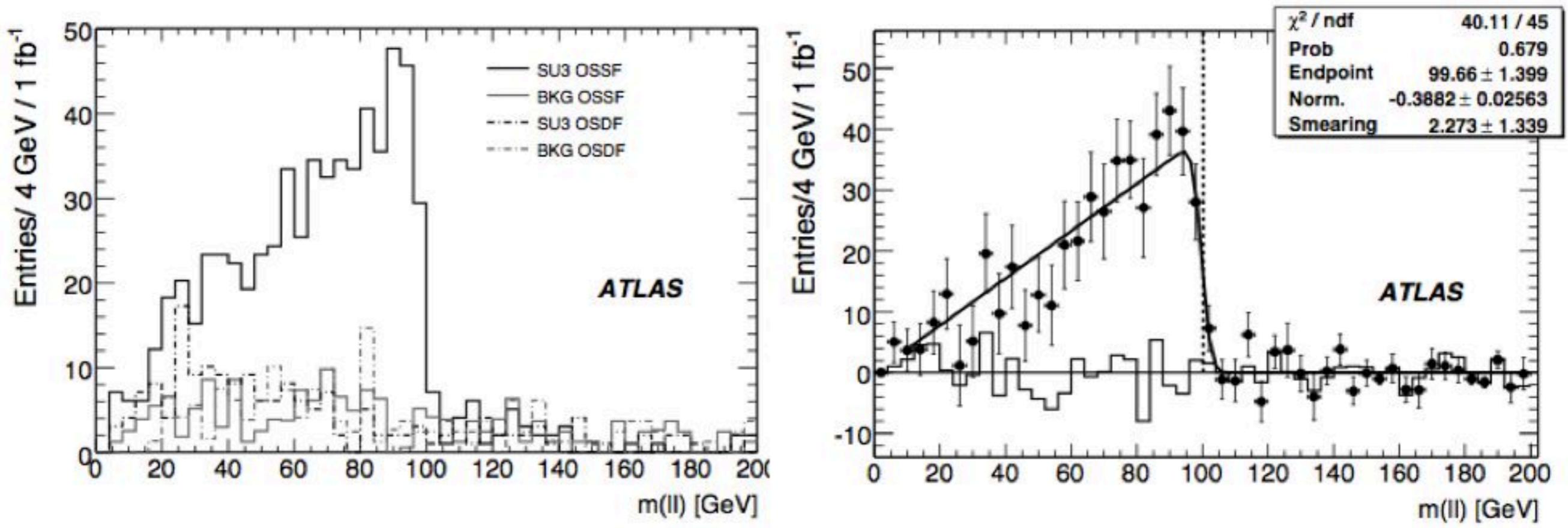
no flavor correlation:
(OSSF)=(OSDF)

$$\text{Flavor subtraction} = (\text{OSSF}) - (\text{OSDF})$$

$$= e^+ e^- + \mu^+ \mu^- - e^+ \mu^- - e^- \mu^+$$

Combinatorics in collider signature: example (II)

$$\tilde{q} \rightarrow \tilde{\chi}_2^0 q \rightarrow l^\pm \tilde{l}^\mp q \rightarrow \tilde{\chi}_1^0 l^\pm l^\mp q$$



taken from ATLAS TDR (2009)

Flavor subtraction = (OSSF) - (OSDF)

$$= e^+e^- + \mu^+\mu^- - e^+\mu^- - e^-\mu^+$$

Combinatorics in collider signature: top-pair

$$t \rightarrow bW^+ \rightarrow bl^+\nu$$

$$\bar{t} \rightarrow \bar{b}W^- \rightarrow \bar{b}l^-\bar{\nu}$$

- * Due to the charge ambiguity on b jets,
of possible pairings = 2:
 $\{l_1, b_1\}$ & $\{l_2, b_2\}$ for right pairing,
 $\{l_1, b_2\}$ & $\{l_2, b_1\}$ for wrong pairing.

$$t_1 + t_2 \rightarrow b_1(p^{b_1})l_1(p^{l_1})\nu_1(k_1) + b_2(p^{b_2})l_2(p^{l_2})\nu_2(k_2)$$

Resolving the combinatorial ambiguity: strategy

- * exploit kinematic variables that have predictable features: edge, threshold, peak, ... → 'test variables'
- * take a partition obeying the features as much as possible (wrong partitions will generally not obey because of no correlation).

Resolving the combinatorial ambiguity: strategy

* exploit **kinematic variables** that have predictable features: edge, threshold, peak, ... → 'test variables'

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only depend on the decay topology
→ 'model-independent'

Combinatorics in collider signature: top-pair

$$(m_{bl}^{(1)})^2 \equiv (p^{b_1} + p^{l_1})^2, \quad (m_{bl}^{(2)})^2 \equiv (p^{b_2} + p^{l_2})^2$$
$$\leq (m_{bl}^{\max})^2 \approx m_t^2 - m_W^2$$

$$(m_{bl}^{(1')})^2 \equiv (p^{b_2} + p^{l_1})^2, \quad (m_{bl}^{(2')})^2 \equiv (p^{b_1} + p^{l_2})^2,$$

no definite cutoff,

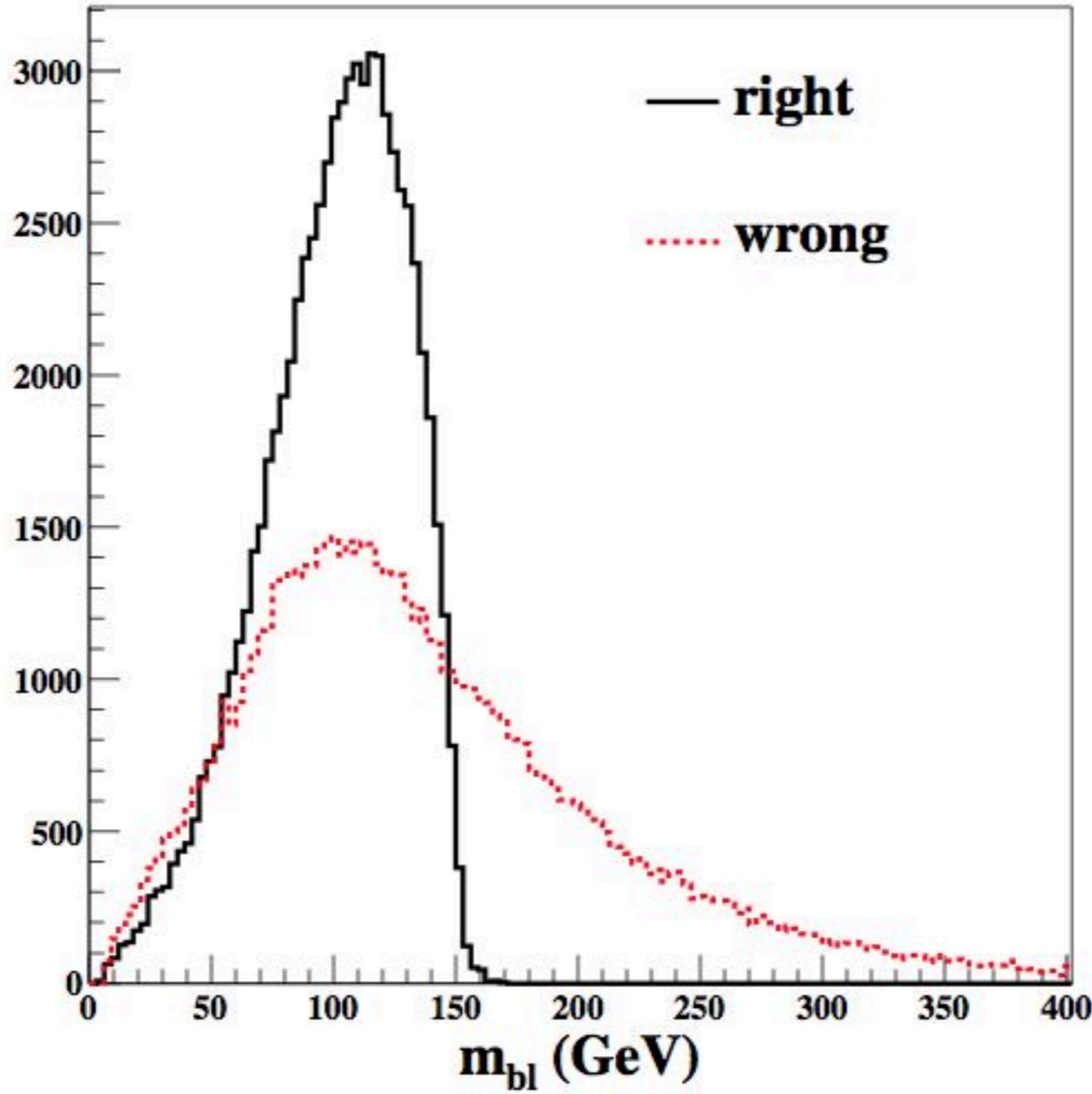
distribution becomes broader for boosted events.

Combinatorics in collider signature: top-pair

LHC, $\sqrt{s} = 7$ TeV

$(m_{bl}^{(1)})^2$
 $\leq (m_{bl}^{max})^2$
 $(m_{bl}^{(1')})^2$
 no definite
 distribution

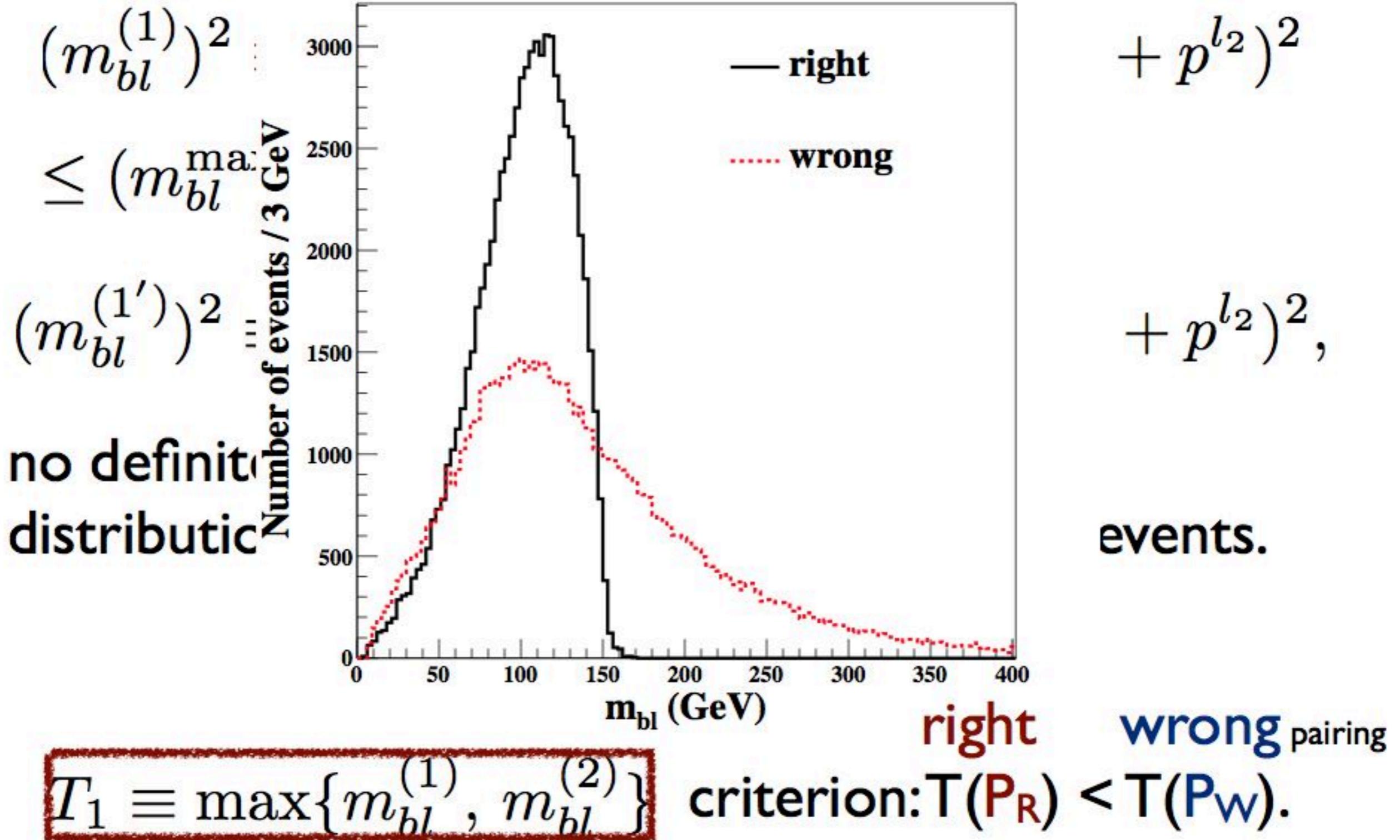
Number of events / 3 GeV



$(m_{bl}^{(2)} + p^{l_2})^2$
 $(m_{bl}^{(2)} + p^{l_2})^2,$
 events.

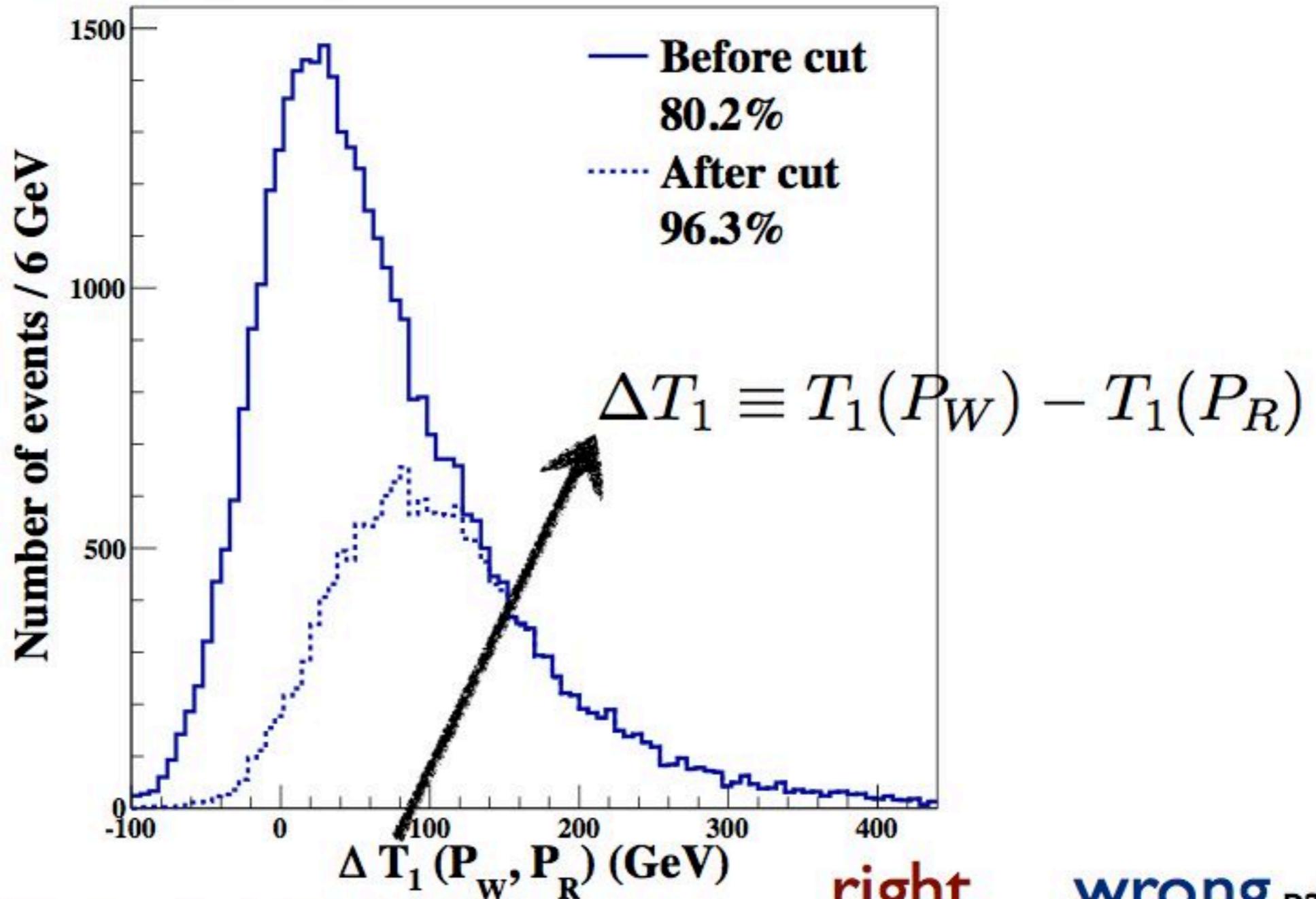
Combinatorics in collider signature: top-pair

LHC, $\sqrt{s} = 7$ TeV



Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV



$$T_1 \equiv \max\{m_{bl}^{(1)}, m_{bl}^{(2)}\}$$

right **wrong** pairing
criterion: $T(P_R) < T(P_W)$.

Combinatorics in collider signature: test variable

C.Lester, D.Summers, PLB(1999)

$$M_{T2} \equiv \min_{\mathbf{k}_{1T} + \mathbf{k}_{2T} = \mathbf{p}_T^{\text{miss}}} \left[\max \left\{ M_T^{(1)}, M_T^{(2)} \right\} \right]$$

$$M_T^{(i)} (m_{bl}^{(i)}, \mathbf{p}_T^{b_i} + \mathbf{p}_T^{l_i}, \mathbf{k}_{iT})$$

: transverse mass in each decay chain

$$M_{T2} \leq m_t \quad \text{for right paring.}$$

Combinatorics in collider signature: test variable

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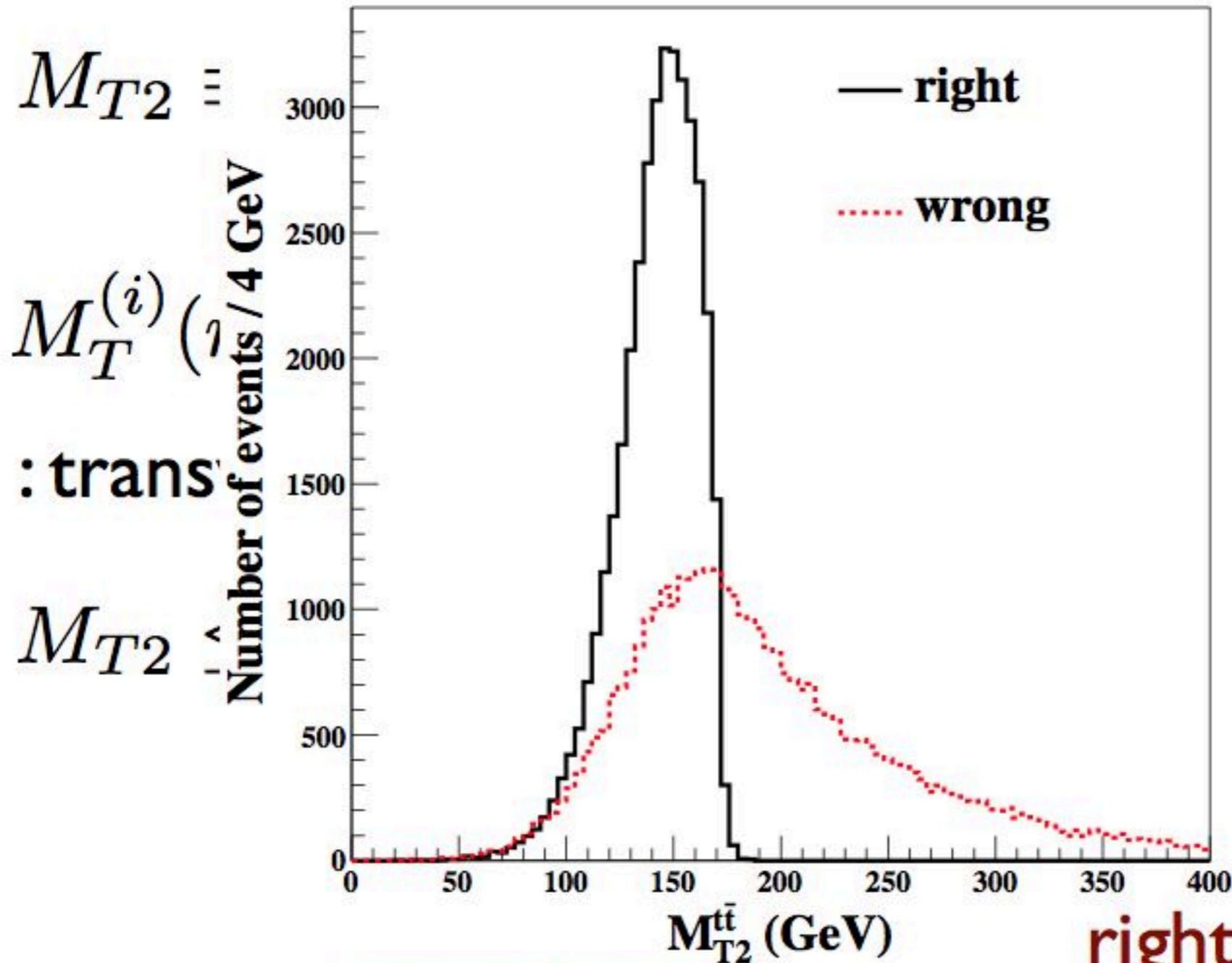
$M_{T2} \leq m_t$ for right paring.

top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

W-pair subsystem (2l+MET): M_{T2}^{WW}

Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV



$\left[T^{(2)} \right]$

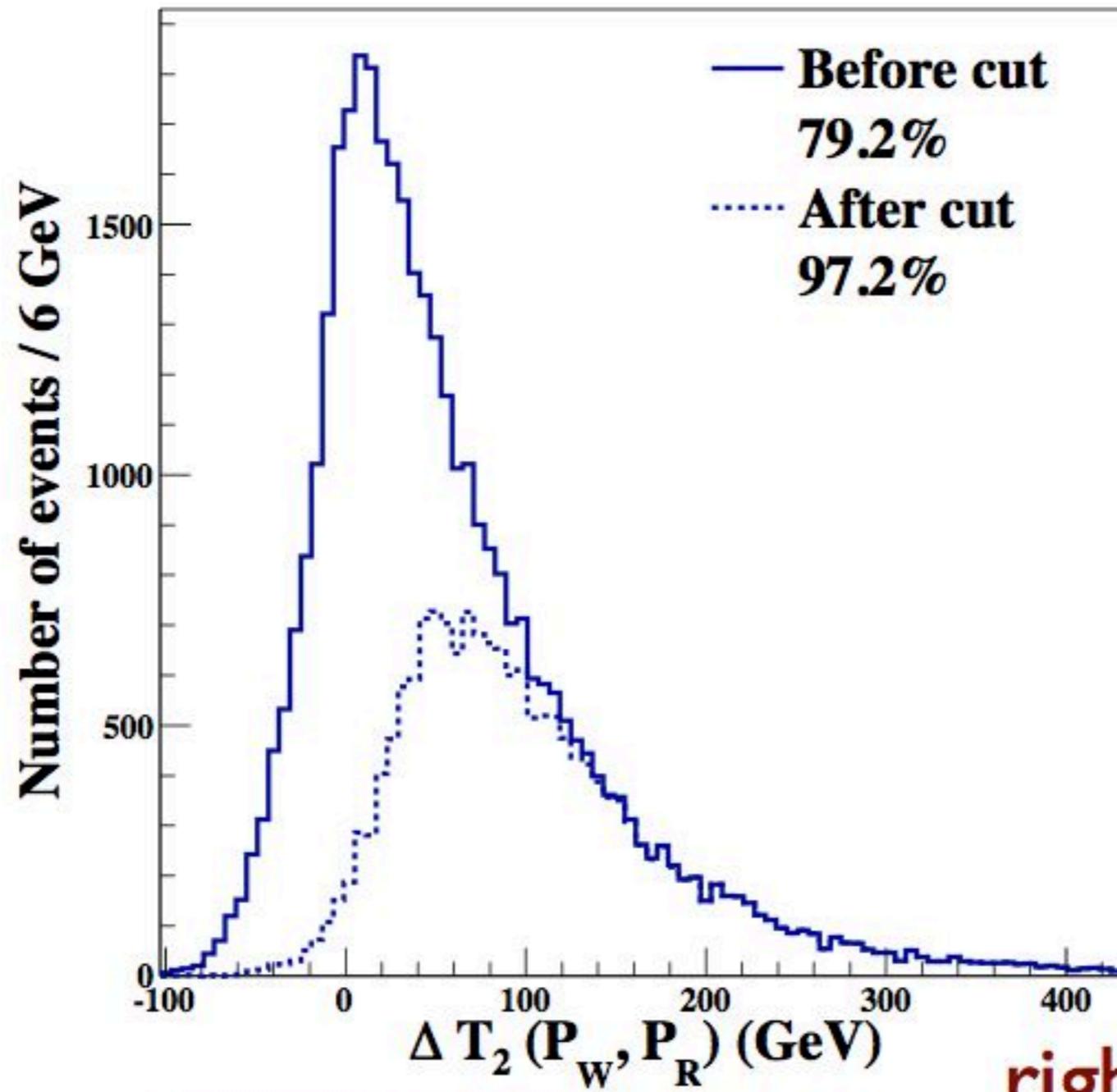
$M_{T2} \equiv$
 $M_T^{(i)}$ (1
 : trans
 $M_{T2} \hat{=}$
 Number of events / 4 GeV

$$T_2 \equiv M_{T2}^{tt}$$

right **wrong** pairing
 criterion: $T(P_R) < T(P_W)$.

Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV



right wrong pairing

$$T_2 \equiv M_{T_2}^{tt}$$

criterion: $T(P_R) < T(P_W)$.

Combinatorics in collider signature: test variable

$$M_{T2} \equiv \min_{\mathbf{k}_{1T} + \mathbf{k}_{2T} = \mathbf{p}_T^{\text{miss}}} \left[\max \left\{ M_T^{(1)}, M_T^{(2)} \right\} \right]$$

also gives the invisible (neutrino) transverse momenta as a result of minimization.

For given k_T values, k_L can be obtained by on-shell eqs.:

$$(p^b + p^l + k)^2 = m_t^2,$$

$$(p^l + k)^2 = m_W^2.$$

top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

W-pair subsystem (2l+MET): M_{T2}^{WW}

Combinatorics in collider signature: test variable

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also gives the invisible (neutrino) momenta

M_{T2} -assisted on-shell (MAOS) reconstruction

... can be obtained by on-shell

$$(p^b + p^l + k)^2 = m_t^2,$$

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top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

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Combinatorics in collider signature: test variable

$$k^{\text{maos}-WW} \quad \text{vs} \quad k^{\text{maos}-t\bar{t}}$$

$$(p^b + p^l + k)^2 = m_t^2,$$

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top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

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Combinatorics in collider signature: test variable

$$k^{\text{maos}-WW} \quad \text{vs} \quad k^{\text{maos}-t\bar{t}}$$

$$\text{If } k^{\text{maos}-WW} = k^{\text{true}}$$

$$\text{then } (p^b + p^l + k^{\text{maos}-WW})^2 = m_t^2$$

$$\text{If } k^{\text{maos}-t\bar{t}} = k^{\text{true}}$$

$$\text{then } (p^l + k^{\text{maos}-t\bar{t}})^2 = m_W^2$$

$$(p^b + p^l + k)^2 = m_t^2,$$

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top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

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Combinatorics in collider signature: test variable

$$k^{\text{maos-WW}}$$

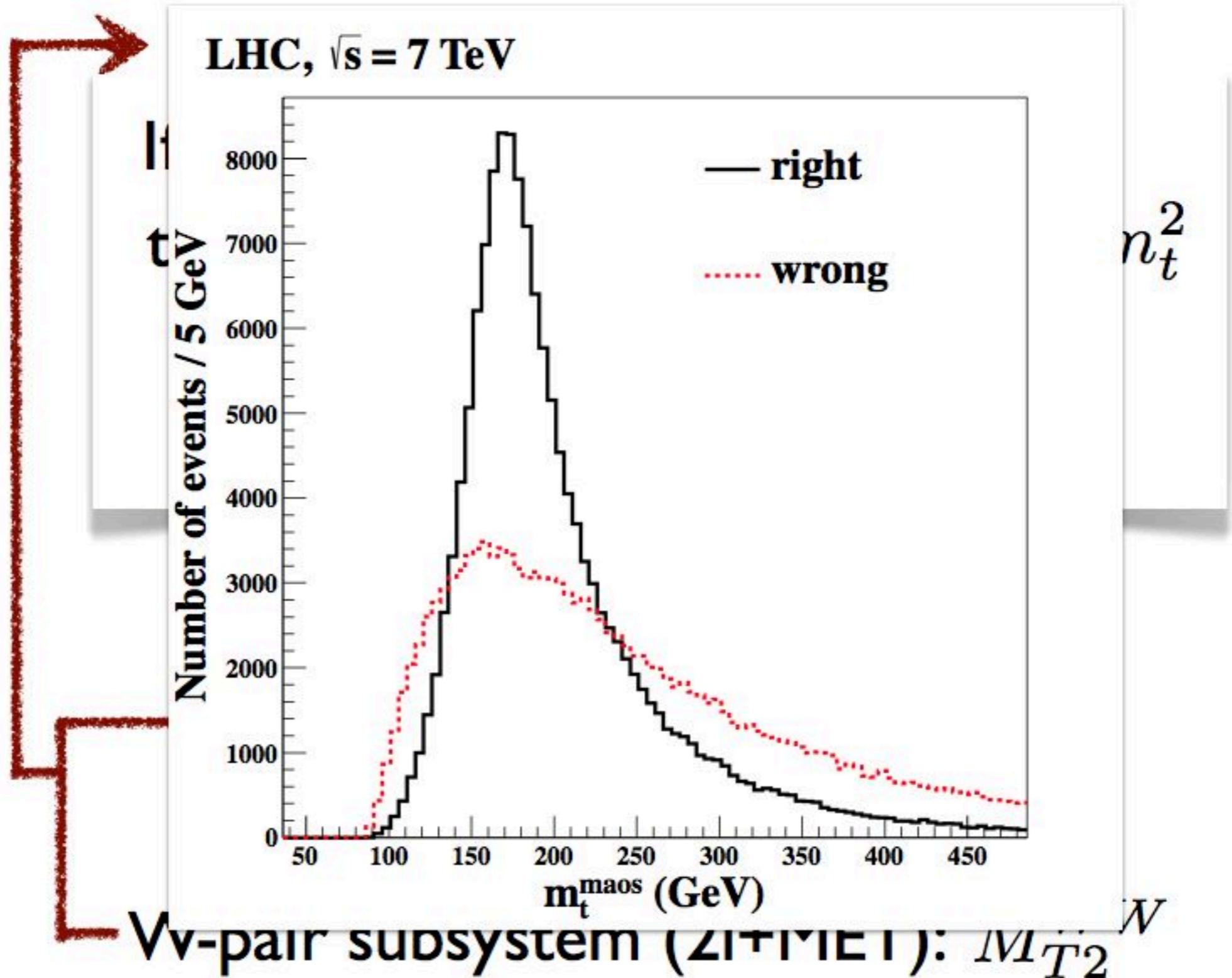
$$\text{If } k^{\text{maos-WW}} = k^{\text{true}}$$

$$\text{then } (p^b + p^l + k^{\text{maos-WW}})^2 = m_t^2 \\ \equiv (m_t^{\text{maos}})^2$$

$$(p^l + k)^2 = m_W^2.$$

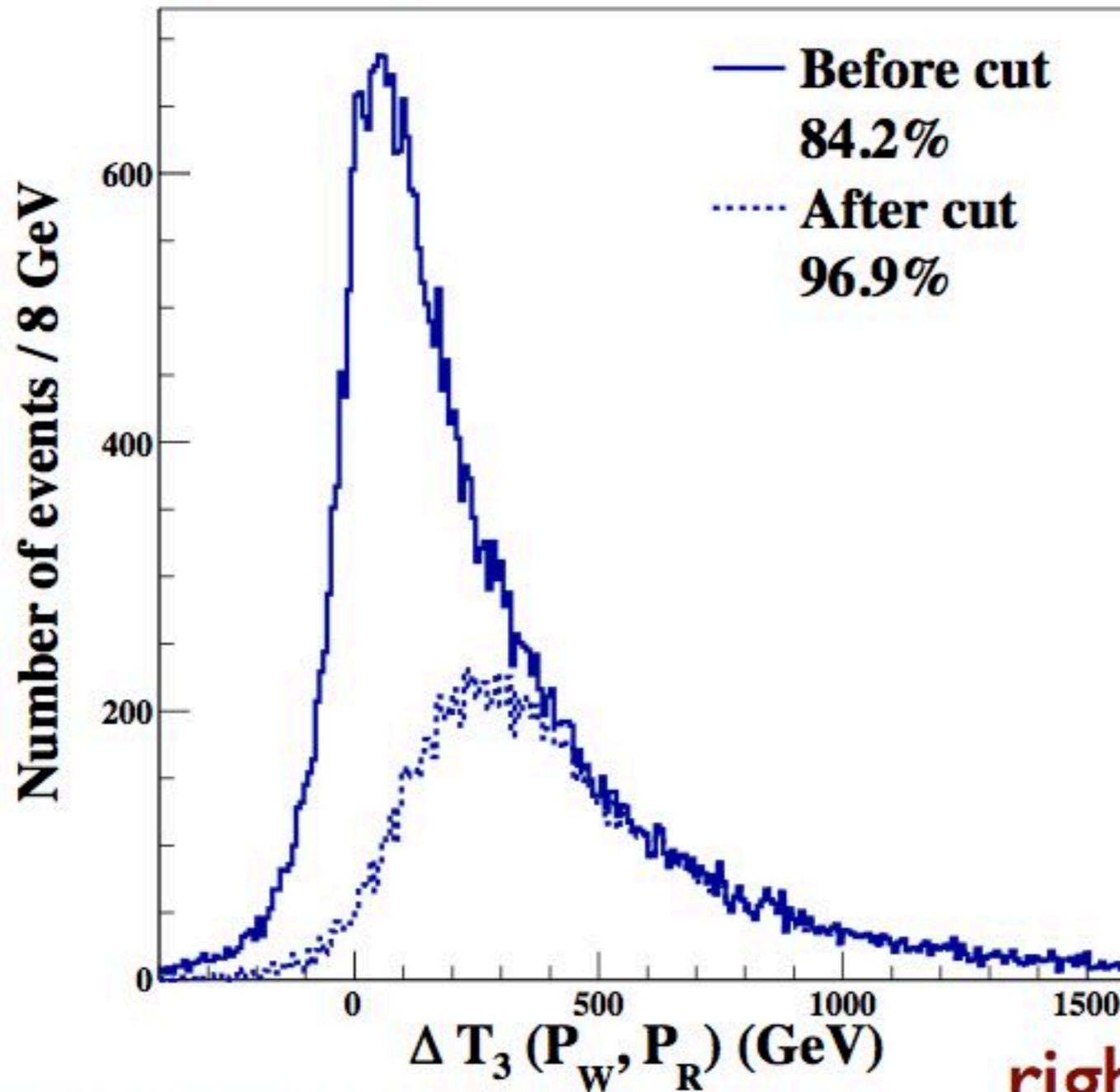
W-pair subsystem (2l+MET): M_{T2}^{WW}

Combinatorics in collider signature: test variable



Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV



$$T_3 \equiv |m_t^{\text{maos}} - m_t|$$

right **wrong** pairing
criterion: $T(P_R) < T(P_W)$.

Combinatorics in collider signature: test variable

$$k^{\text{maos-}t\bar{t}}$$


If $k^{\text{maos-}t\bar{t}} = k^{\text{true}}$
then $(p^l + k^{\text{maos-}t\bar{t}})^2 \equiv (m_W^{\text{maos}})^2$
 $= m_W^2$

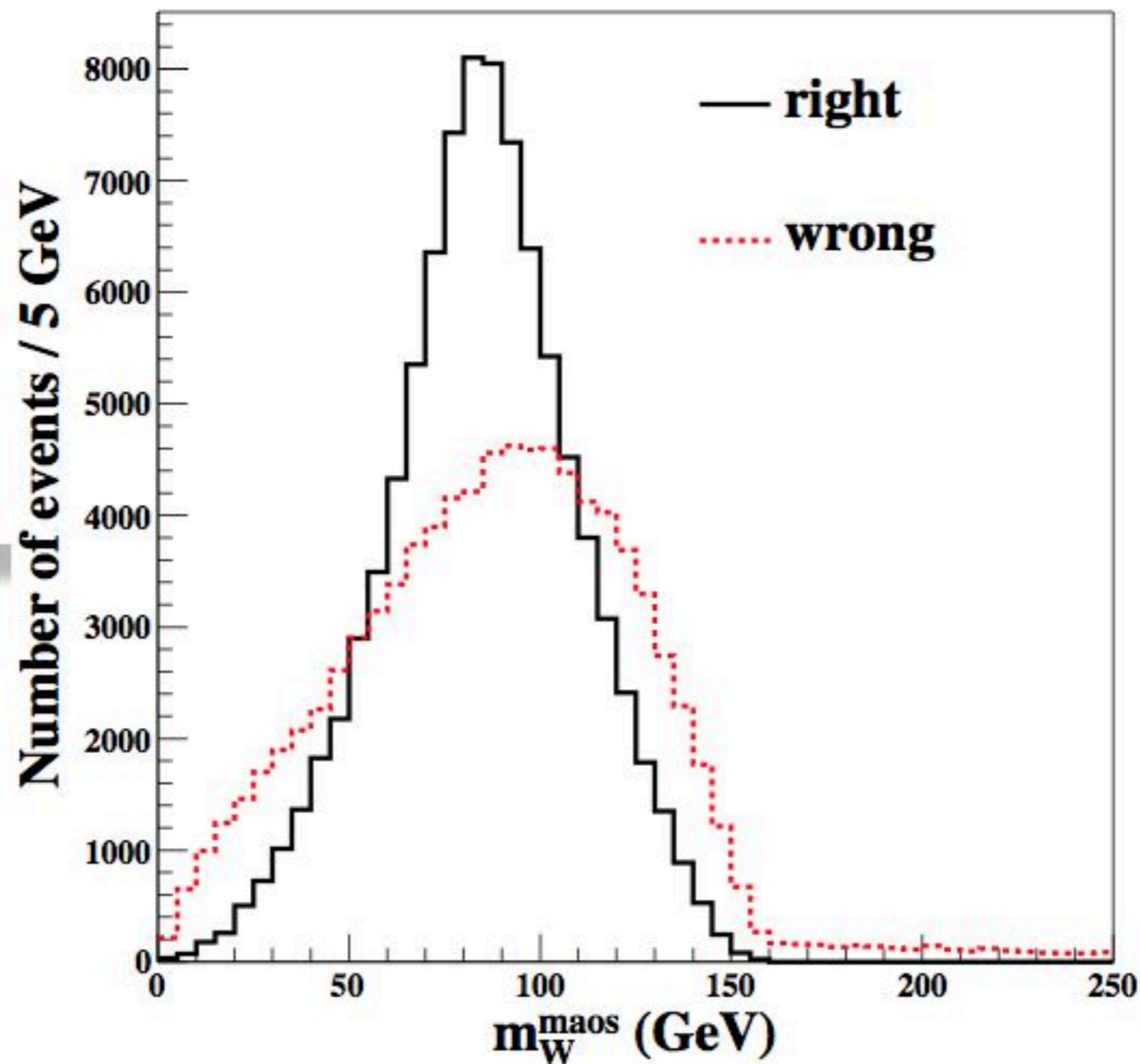
$$(p^b + p^l + k)^2 = m_t^2,$$


top-pair full system (2b+2l+MET): $M_{T2}^{t\bar{t}}$

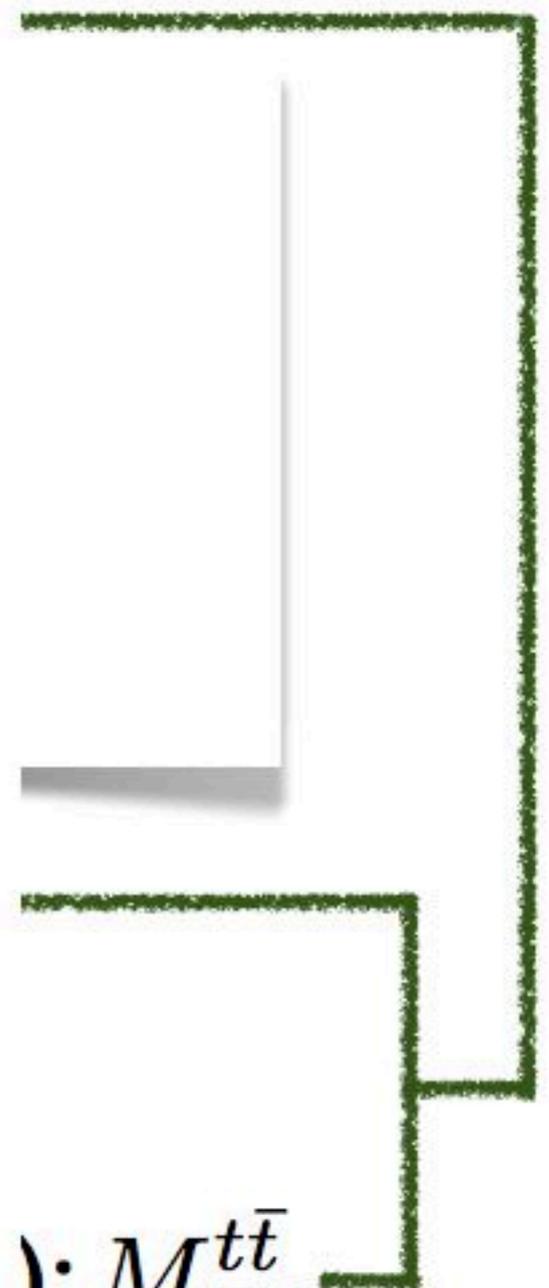


Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV

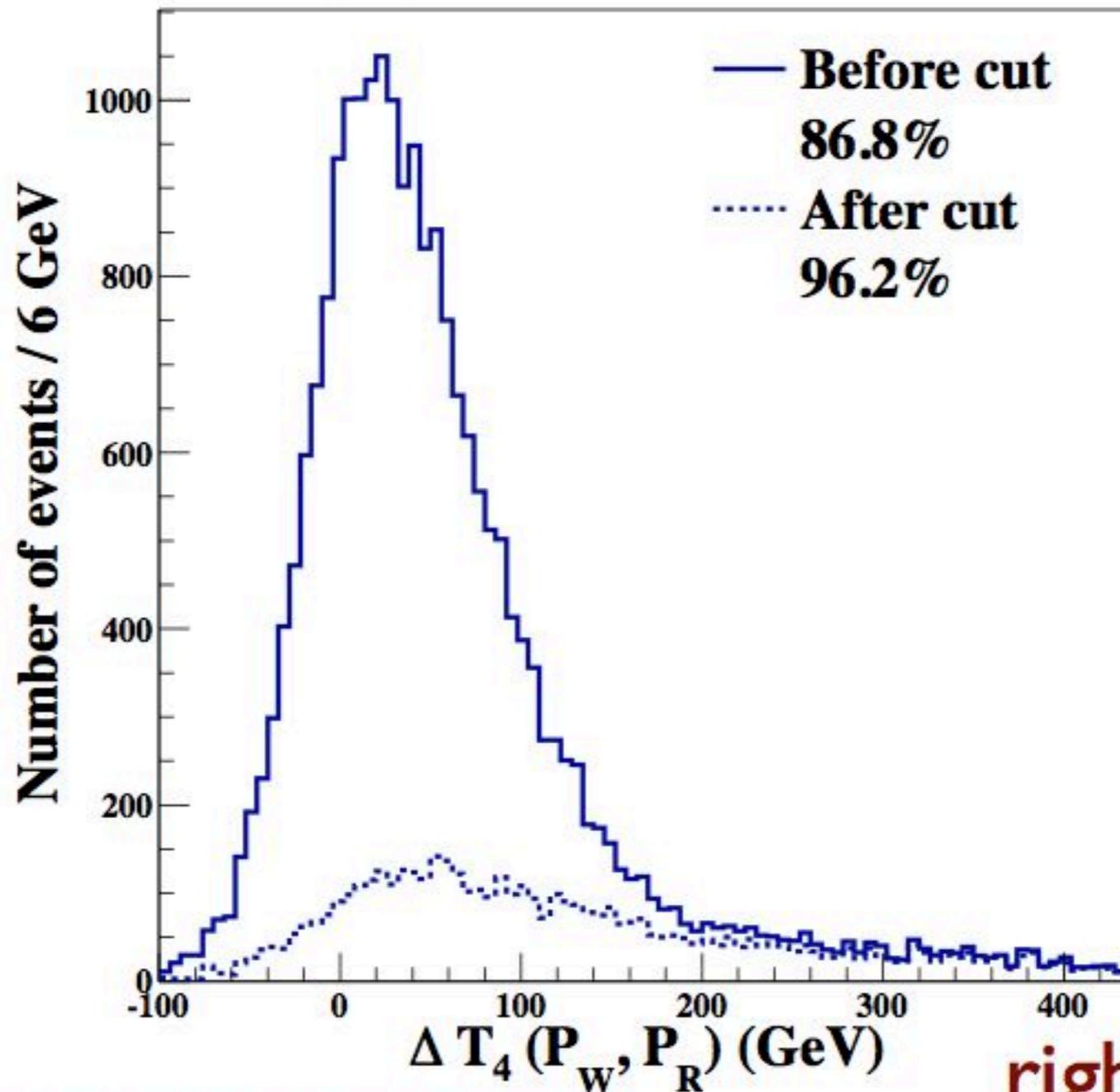


) : $M_{T2}^{t\bar{t}}$



Combinatorics in collider signature: test variable

LHC, $\sqrt{s} = 7$ TeV



$$T_4 \equiv |m_W^{\text{maos}} - m_W|$$

right **wrong** pairing
criterion: $T(P_R) < T(P_W)$.

Combinatorics in collider signature: test variables

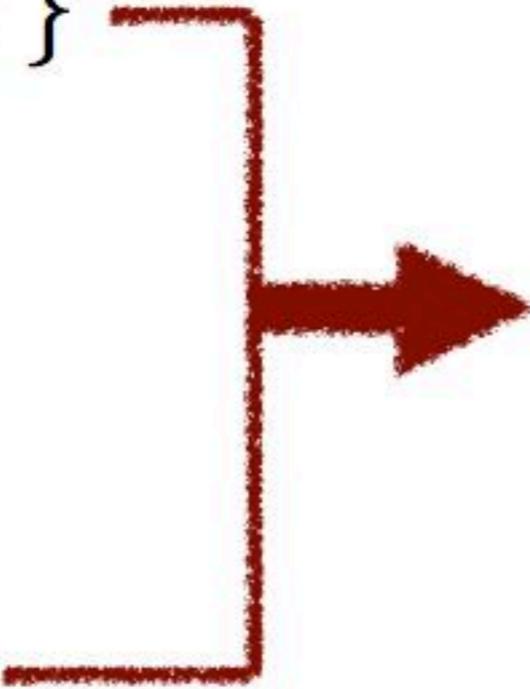
$$t_1 + t_2 \rightarrow b_1(p^{b_1})l_1(p^{l_1})\nu_1(k_1) + b_2(p^{b_2})l_2(p^{l_2})\nu_2(k_2)$$

$$T_1 \equiv \max\{m_{bl}^{(1)}, m_{bl}^{(2)}\}$$

$$T_2 \equiv M_{T_2}^{t\bar{t}}$$

$$T_3 \equiv |m_t^{\text{maos}} - m_t|$$

$$T_4 \equiv |m_W^{\text{maos}} - m_W|$$



combined method?

* take a partition obeying the features as much as possible.

Combinatorics in collider signature: test variables

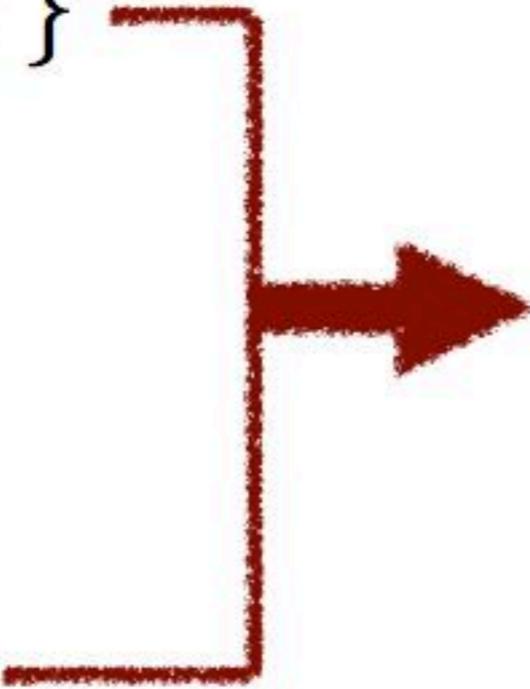
$$t_1 + t_2 \rightarrow b_1(p^{b_1})l_1(p^{l_1})\nu_1(k_1) + b_2(p^{b_2})l_2(p^{l_2})\nu_2(k_2)$$

$$T_1 \equiv \max\{m_{bl}^{(1)}, m_{bl}^{(2)}\}$$

$$T_2 \equiv M_{T_2}^{t\bar{t}}$$

$$T_3 \equiv |m_t^{\text{maos}} - m_t|$$

$$T_4 \equiv |m_W^{\text{maos}} - m_W|$$



combined method?

* take a partition obeying the features as much as possible.

→ select a partition with the smallest value of T_i for the largest number of test variables.

(~ **89%** efficiency without loss of statistics)

Combinatorics in collider signature: test variables
+ event selection cut

wrong partitions: no definite cutoff,
distribution becomes broader for boosted events.

Combinatorics in collider signature: test variables

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wrong partitions: no definite cutoff,
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Partition-insensitive kinematic variables:

$M_T^{t\bar{t}}$ (transverse mass of top-pair full system)

m_V (invariant mass of $2b+2l$ system)

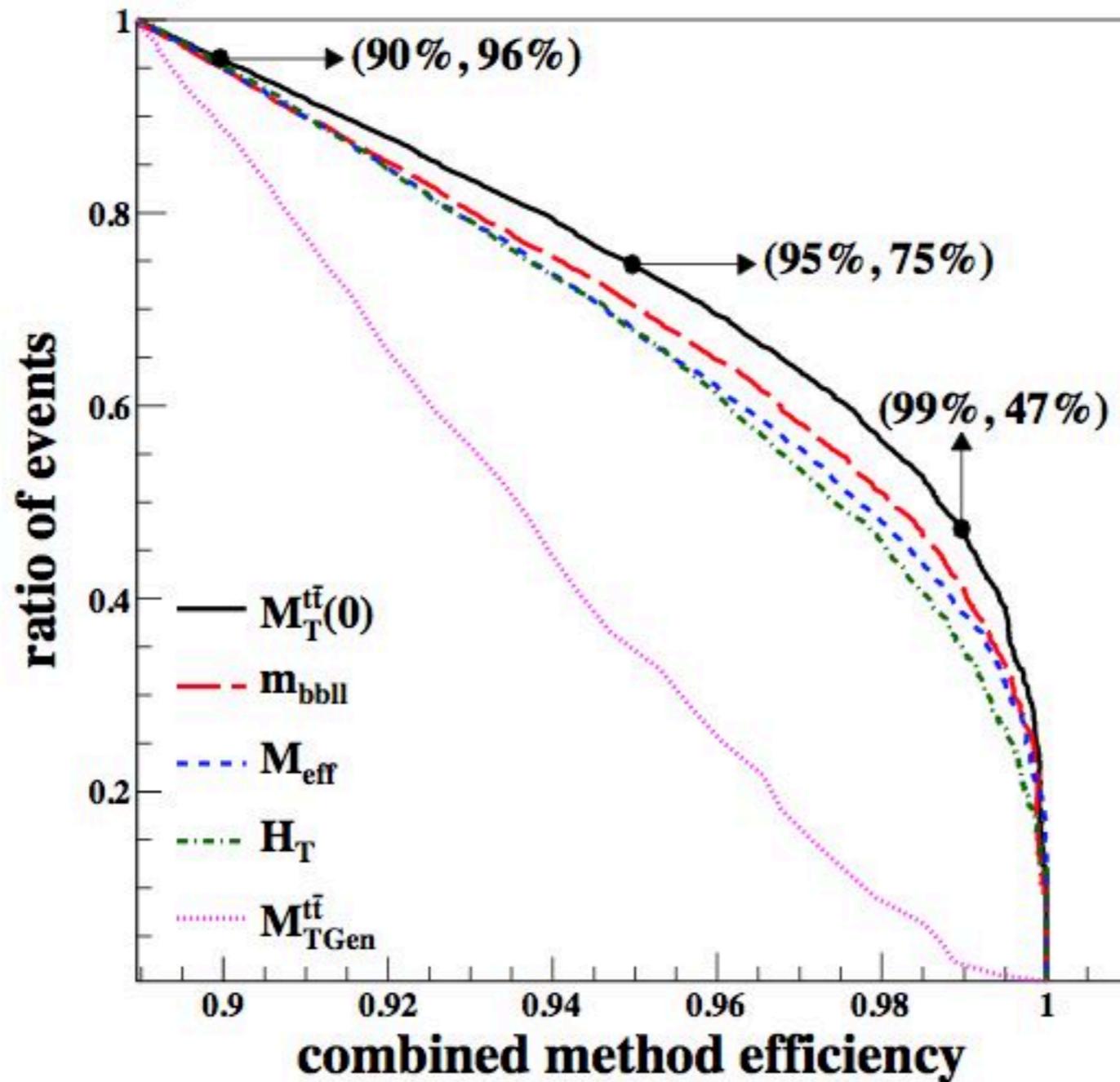
M_{eff} (scalar sum of all P_T + missing E_T)

H_T (scalar sum of all P_T)

$M_{T\text{Gen}}^{t\bar{t}}$ (smallest M_{T2} of all possible in full system)

Combinatorics in collider signature: test variables + event selection cut

LHC, $\sqrt{s} = 7$ TeV



Summary and Outlook

- * We proposed a novel method to resolve combinatorial ambiguities in hadron collider events involving two invisible particles in the final state.
- * The method is based on the kinematic variables with the assumption of decay topology.
- * It can be utilized for new physics processes with pair production and WIMP - in case of pairing and/or ordering ambiguities - (e.g. R-parity conserving SUSY), as well as the dileptonic SM top-pair process.
- * For practical use, we will also study the effects of momentum smearing, ISR, and poor mass information.